



Accuracy of TEXTOR He-beam diagnostics

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Introduction – standard He-beam diagnostics on TEXTOR

Atomic processes involved

Evaluation procedures

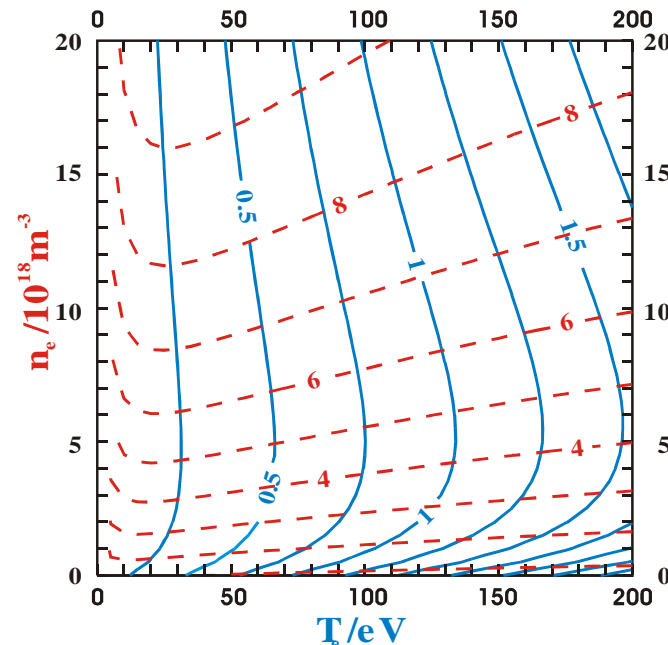
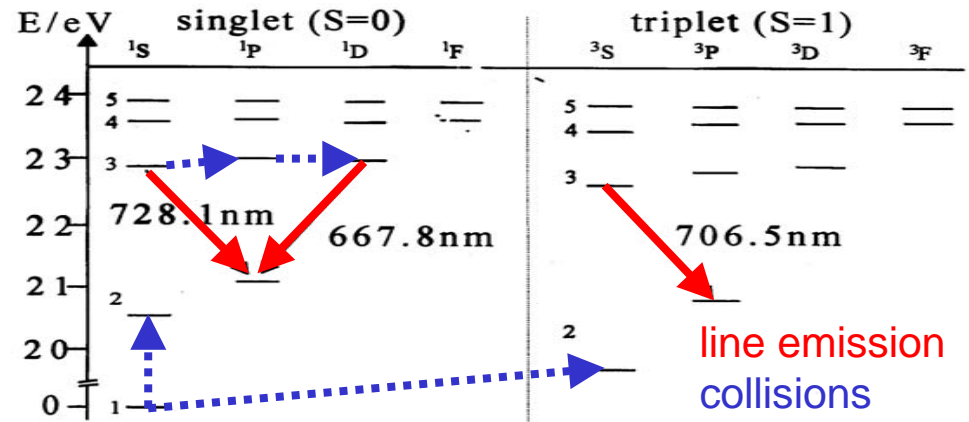
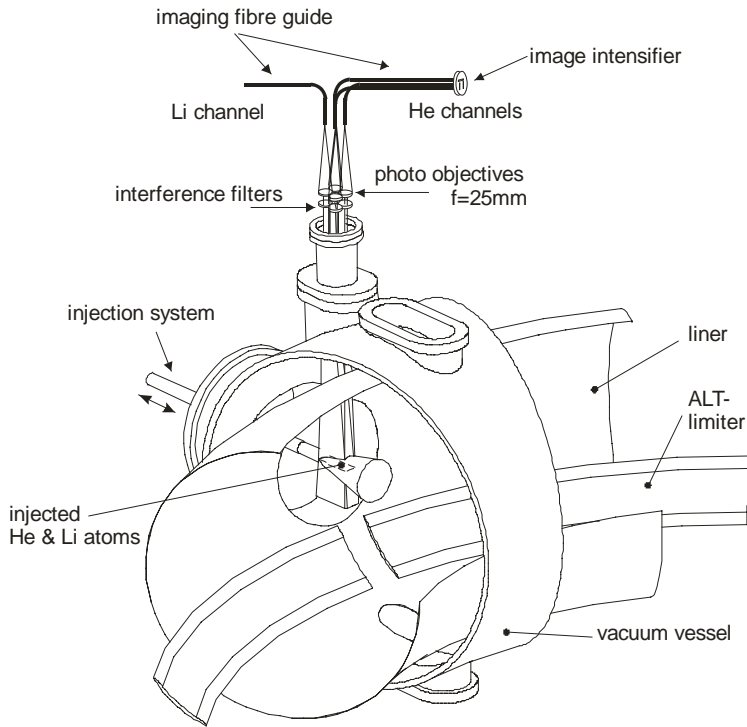
Atomic data - standard, new

Results - comparison of intensity profile fits
 - comparison for different atomic data
 - comparison for different discharge conditions
 - comparison of different diagnostics

Summary & conclusions



Standard TEXTOR He-beam diagnostics



line intensity ratios:
singlet : singlet
 $I(668\text{nm}) / I(728\text{nm})$
 n_e dependent - - -

singlet : triplet
 $I(728\text{nm}) / I(706\text{nm})$
 T_e dependent —

M.Brix (2000)

- $n_i/n_1 \cong \text{constant}$
- relaxation time τ_r sufficiently small
- temporal resolution: $\Delta t = \tau_r(\text{max})$
- spatial resolution: $\Delta x = \Delta t \cdot v_{\text{Strahl}}$
- no p & d collisions

Atomic processes

$$v_b \frac{d}{dx} n_i =$$

pop. processes

$$+ \sum_{j,j \neq i} \langle \sigma_{e,ji} v \rangle n_e n_j$$

$$+ \sum_{j,j \neq i} \langle \sigma_{p,ji} v \rangle n_{ion} n_j$$

$$+ \sum_{j,j > i} A_{ji} n_j$$

depop. processes

$$- \sum_{i,j \neq i} \langle \sigma_{e,ij} v \rangle n_e n_i$$

$$- \sum_{i,j \neq i} \langle \sigma_{p,ij} v \rangle n_{ion} n_i$$

$$- \sum_{j,j < i} A_{ij} n_i$$

$$- \langle \sigma_{e \rightarrow ionis,i} v \rangle n_e n_i$$

$$- \langle \sigma_{p \rightarrow ionis,i} v \rangle n_{ion} n_i$$

$$- \langle \sigma_{cx,i} v \rangle n_{ion} n_i$$

electron collision
excitation

ion collision
excitation

spontaneous
emission

electron collision
ionisation

ion collision
ionisation

charge
exchange

Evaluation procedures

For known $T_e(r)$ and $n_e(r)$ -> line intensity profiles $I_\lambda(r) = n_k A_{ki}$

However, we need for known $I_\lambda(r) = n_k A_{ki}$ -> $T_e(r)$ and $n_e(r)$

M. Brosda and M. Brix have already shown that it is possible to restore T_e and n_e profiles from line profiles $I_1(r); I_2(r); I_3(r)$

Possible approaches:

non-stationary (NS): direct solution of the general (time / space) dependent equation and fit $I_\lambda(r) = n_k A_{ki}$ to the measured one.

quasi-stationary (QS): approx. solution of the NS equations: on the right the St-solution of the balance equations for rel. populations n_k/n_1 of the excited states is inserted.

stationary (St): the derivatives are neglected assuming that the time constant of the measured processes are larger than the relaxation times τ for the transitions used.

Atomic data - levels

The following atomic levels of He I have been included in the model:

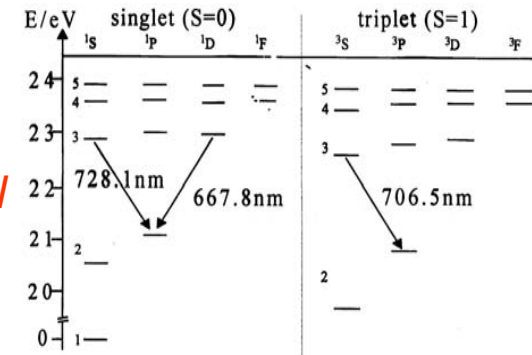
(1): $1s^2\ ^1S$; $1snl\ ^1L$; $1snl\ ^3L$; $n = 2; 3; 4$, $l = 0; 1; 2$

(1i): $1s4f\ ^1F$; 3F

(2): $1sn_{-1}$ (singlets), $1sn_{-3}$ (triplets), $n = 5; 6; 7$ summed over all l

(3): $1sn$; $n = 8; 9$ summed over all l and S

(c): $c = 1s$ ground state of He II



"effective" levels have been introduced, which describe group of levels summed over some quantum numbers (decreases drastically the dimensions of the statistical matrix).

group (1): SL coupling is a good approximation.

group (1i) and any levels with $l > 2$: deviation from the SL coupling is significant.

The matrix of the eigen-vectors in intermediate coupling was obtained for the states (1i) using the GRASP92-code (*Drake*).

For the levels $1sn$; $n = 8; 9$ summed over all l and S the type of coupling is not important.

For the levels $1sn_{-1}$ (singlets), $1sn_{-3}$ (triplets), $n = 5; 6; 7$ an effective mixing due to deviation from the SL coupling was introduced.

Atomic data – energies, radiation, collision with e

The following atomic levels of He I have been included in the mode

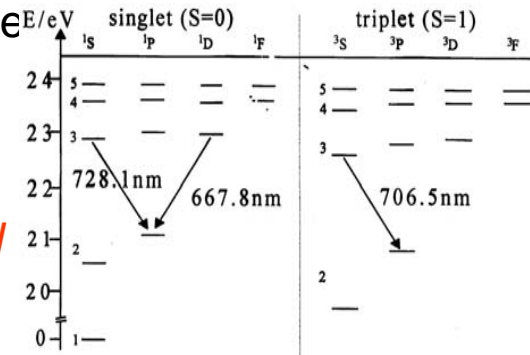
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Energy levels: from NIST

A for $1s2p\ ^{1,3}P - 1s^2\ ^1S$ from *Wiese* (NSRDS-N.B.S.)

A for groups (1), (1i) from ATOM, intermediate coupling for $1s4f - 1s3d$, SL for others

A for transitions inside and between groups (2), (3): Kramers formula

$\langle\sigma v\rangle_{\text{ex}}$ for group (1): CCC-89 cross sections (*Bray*)

$\langle\sigma v\rangle_{\text{ex}}$ from group (1) to groups (1i),(2),(3),c: ATOM code (Norm.BA)

$\langle\sigma v\rangle_{\text{ex}}$ inside groups (1),(2),(3): semiclassical method (*Beigman*)

transitions to group (1i): intermediate coupling

Atomic data – collisions with p,d, CEX

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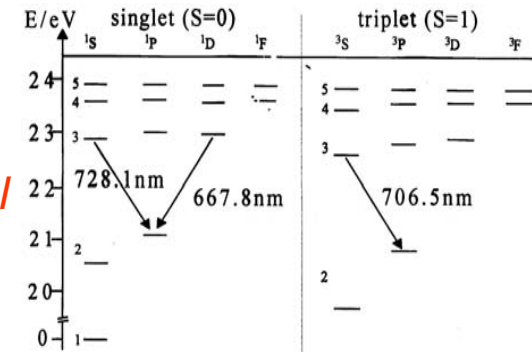
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Collisions with heavy particles at thermal energies may only contribute for transitions with $\Delta n = 0$

$\langle \sigma v \rangle_{p,d}$ for the most important $3l - 3l'$: from CC-method (code ATCC (*Borodin*))
for all other transitions: Norm. Born approximation (*Borodin*)

$\langle \sigma v \rangle_{CEX}$ from the metastable state $1s2s$ from code CAPT (*Shevelko*) with cross sections for $n \geq 4$ from (*Janev*)

Normalized Born Approximation

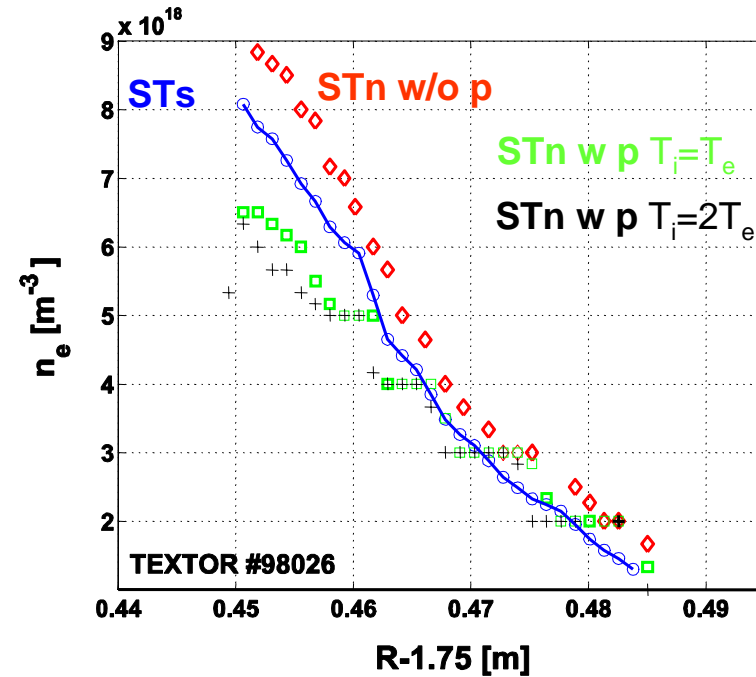
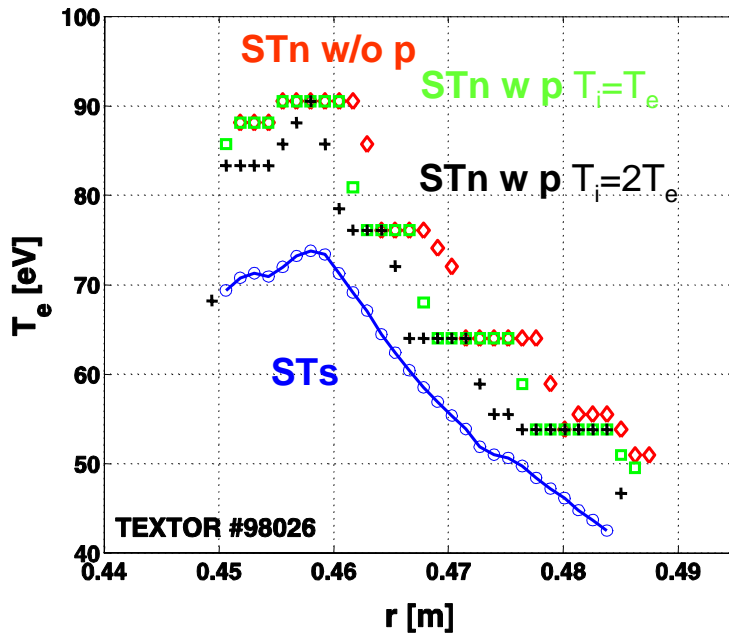
$$\langle v \sigma_p^B \rangle = z^{-3} \frac{A \beta_\mu^{1/2}}{(\beta_\mu + \chi)(\beta_\mu / b_D + 1)} \cdot \exp(-\Delta E / kT);$$

$$\beta_\mu = Z^2 R y / kT \cdot (\mu / m); \quad b_D = 2^{5D-6};$$

Close coupling method

$$\langle v \sigma_p^C \rangle = \langle v \sigma_p^B \rangle \cdot \exp(-\sqrt{(D/50) \beta_\mu})$$

Results – Comparison of different models



proton collisions do not seem to play a strong role

some effect is of course expected at high densities – T_i effects are still weak

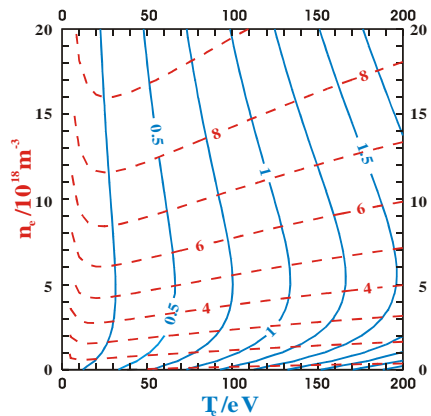
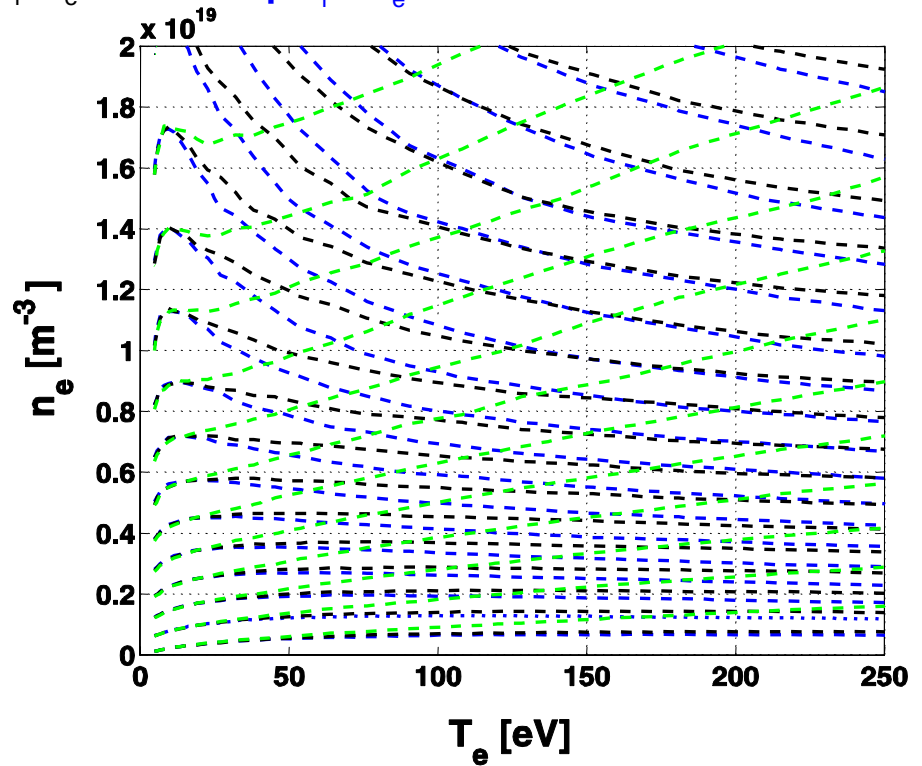
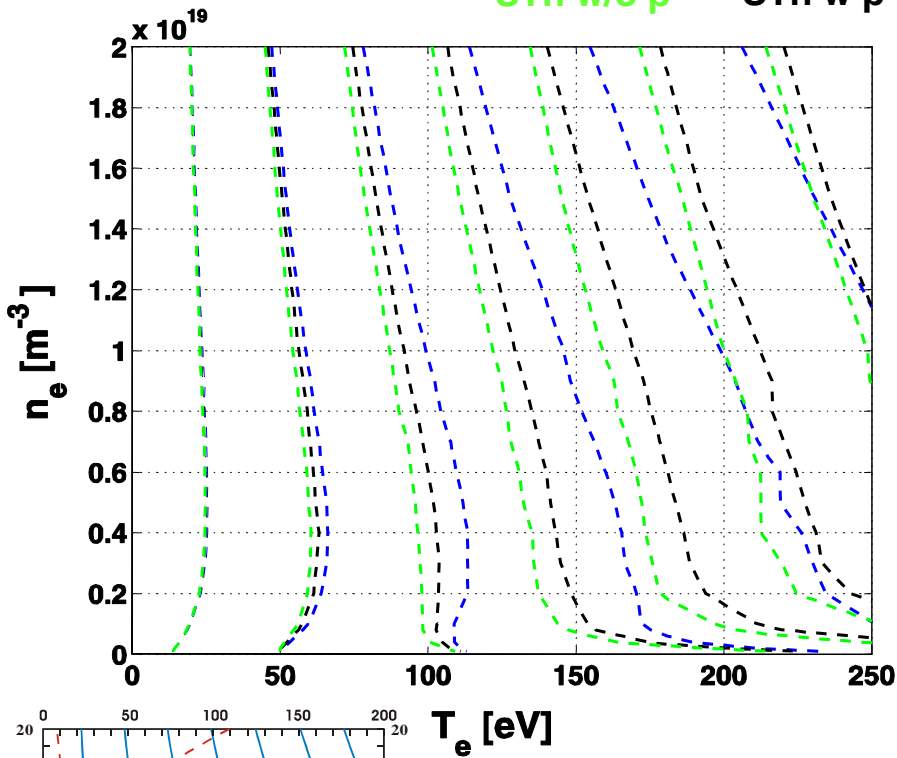


Results – T_e & n_e diagram (new)

STn w/o p

STn w p $T_i=T_e$

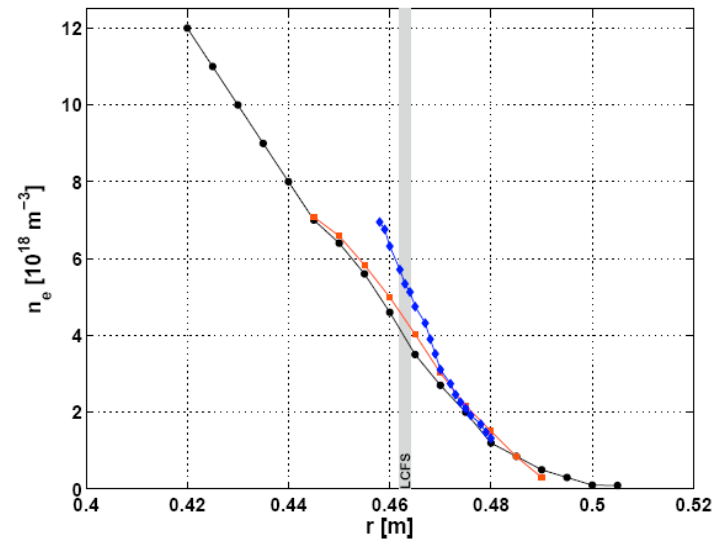
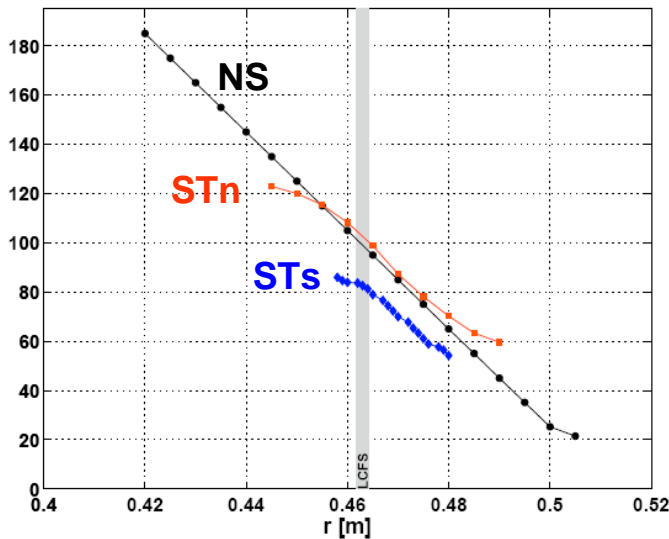
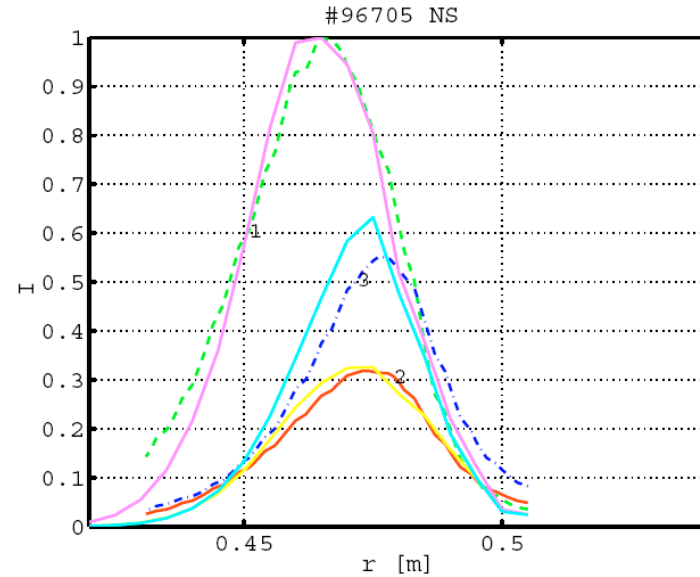
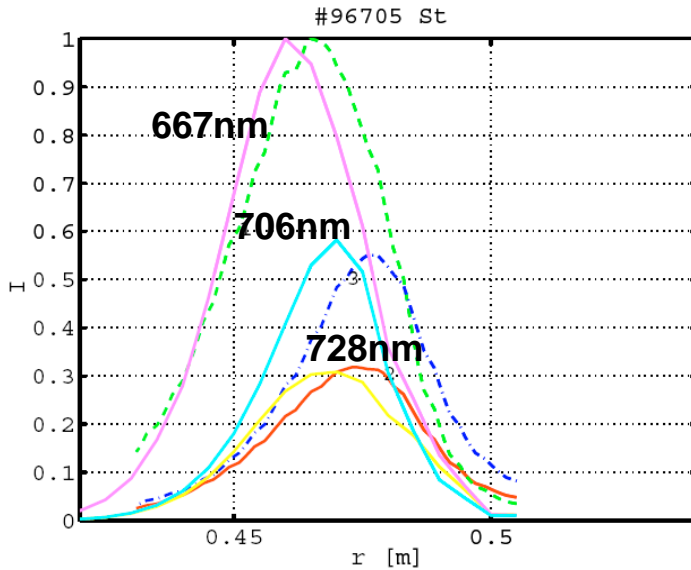
STn w p $T_i=2T_e$



evaluation ranges could be extended
 proton collisions reduce the electron densities
 " " increase the electron temperatures



Results – Comparison for intensity profile fits (Ohmic discharge) (# 96705)



for large radii:

τ_r is too large

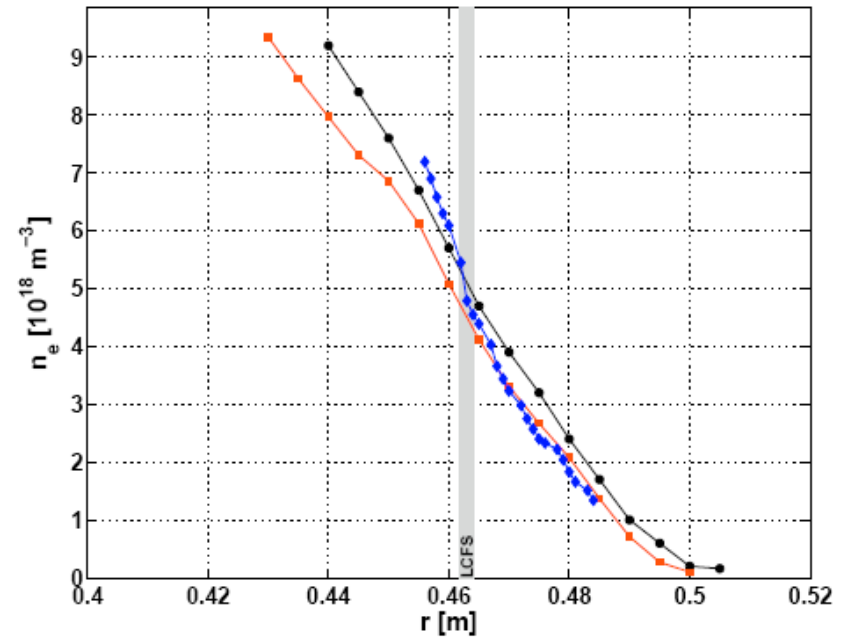
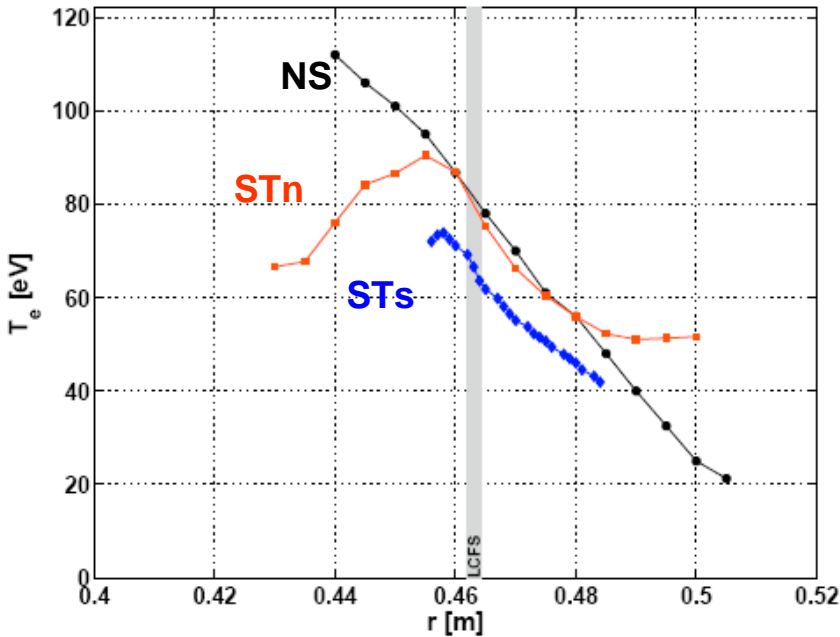
for small radii:

stronger impact of higher I-mixing

(not yet considered)



Results – Comparison for Ohmic discharges (# 98026)



for large radii:

τ_r is too large

for small radii:

stronger impact of
higher I-mixing

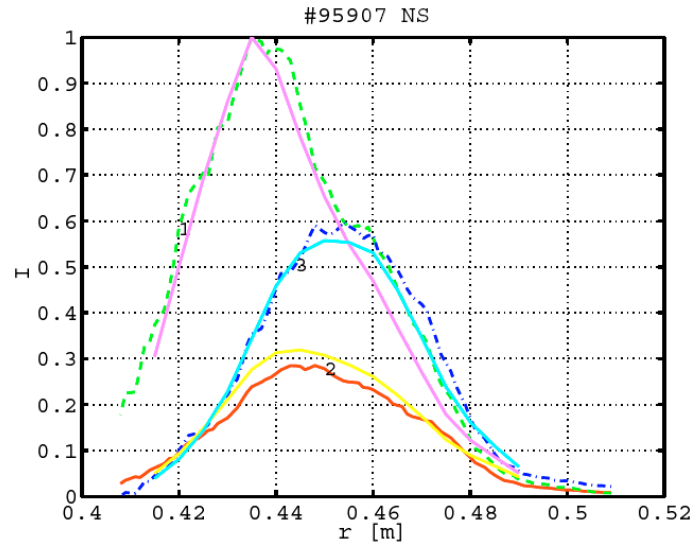
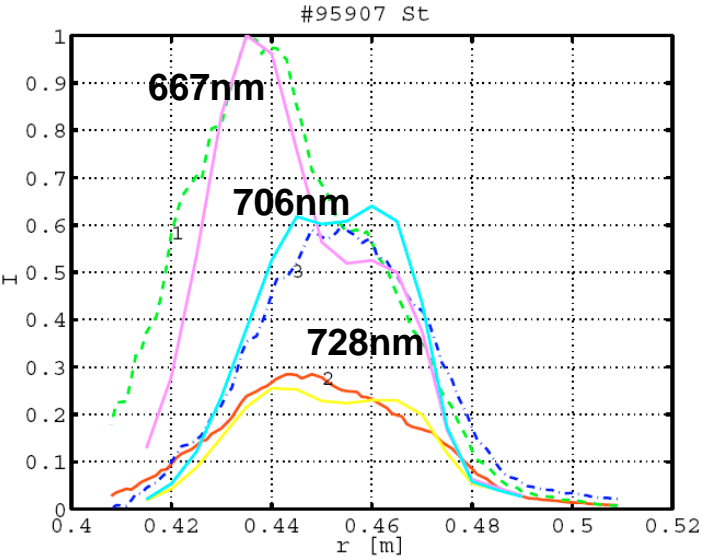
(not yet considered)

For Ohmic discharges:

The deviation between Brix's CRM and the new one accounts to 10% for T_e and is negligible for n_e .

The impact of proton collisions as well as of the new atomic processes (CXRS and a new ionisation rate coefficient) seems to have a minor effect.

Results – Comparison for intensity profile fits (small NBI-heating) (# 95907, 300kW)

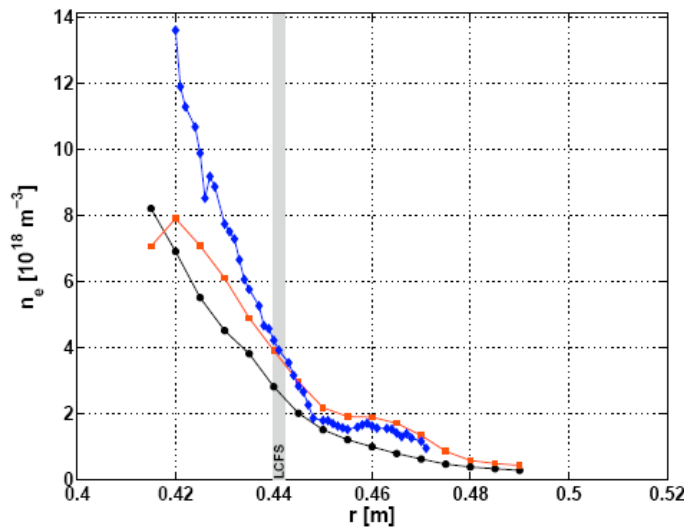
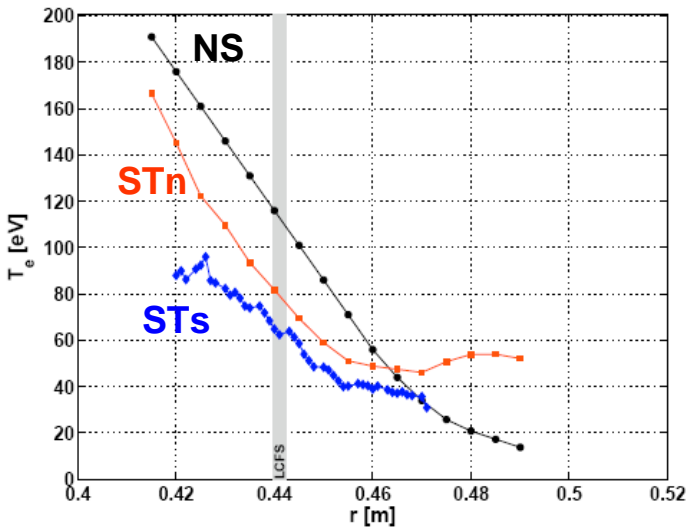


T_e & n_e

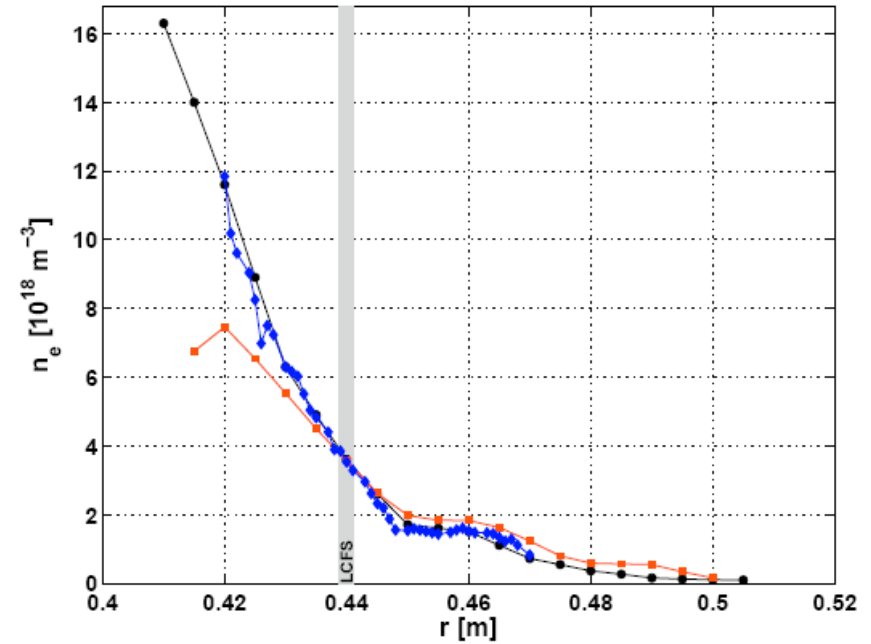
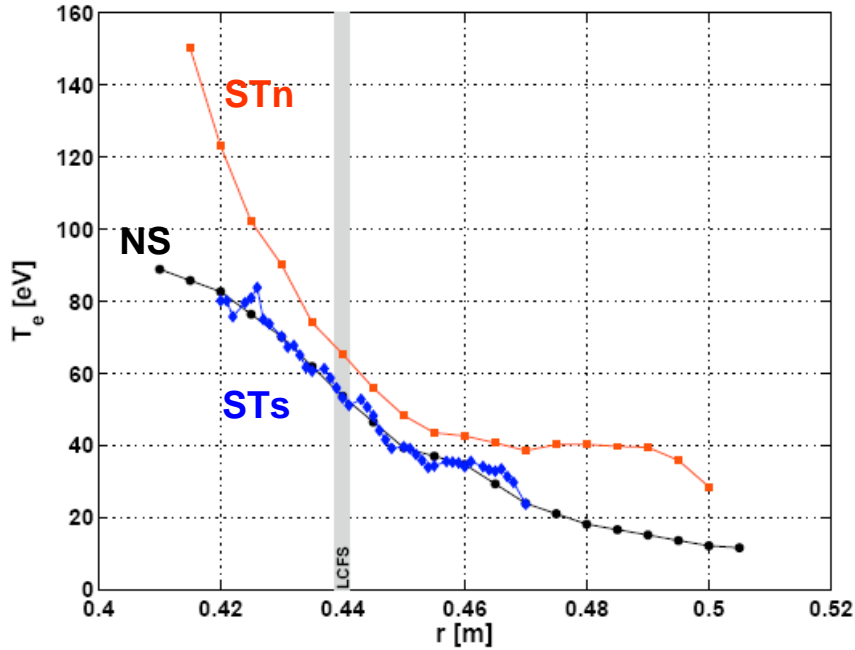
strong deviations (especially for T_e)

for all models at small radii

Influence from higher levels !?



Results – discharges with small NB-heating (# 95896, 300kW)

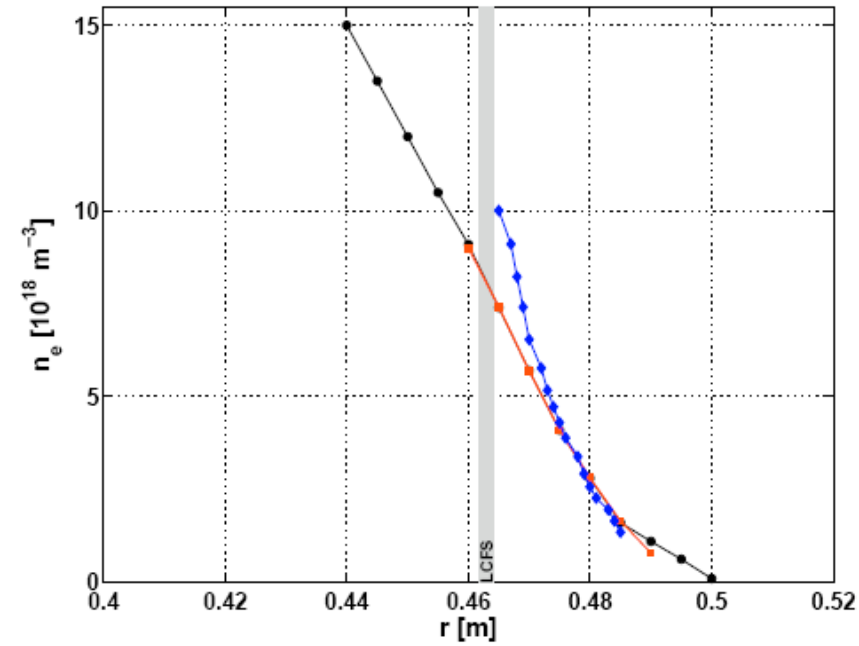
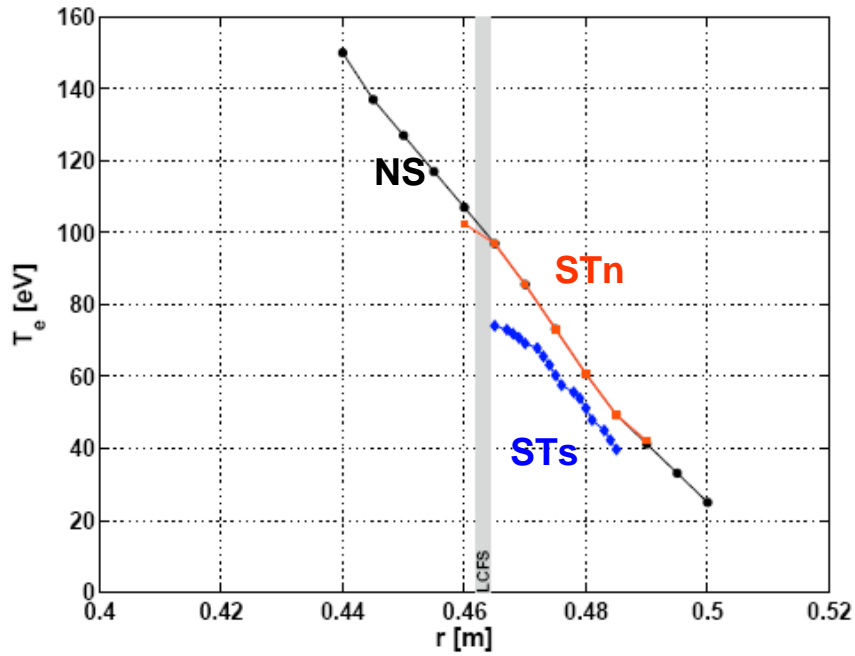


T_e & n_e

much better coincidence now



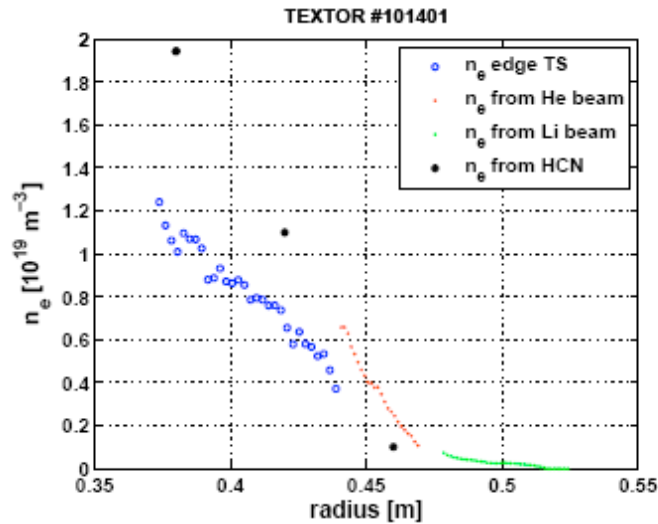
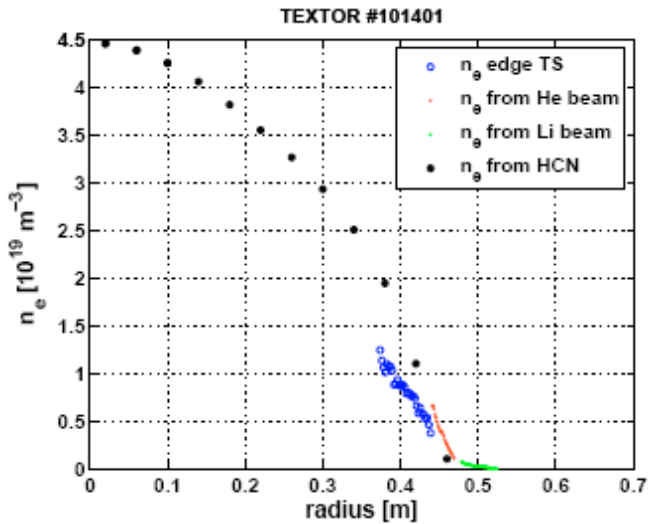
Results – discharges with strong NB-heating (# 96710, 1200kW)



T_e & n_e

again much better coincidence now, **NS** allows extension into the confined plasma

Results – comparison with different diagnostics

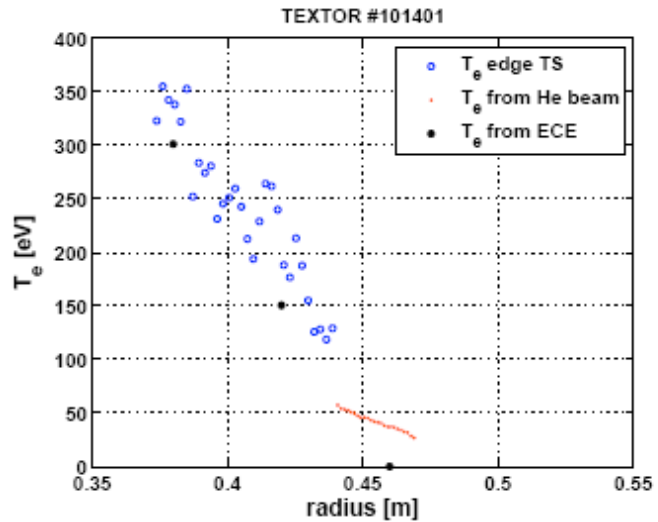
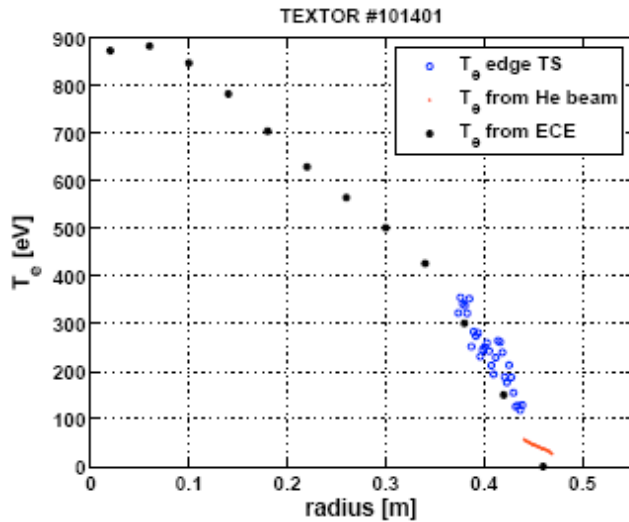


n_e

He- & Li-beam data extrapolate well into the HCN data

TS seems too small (calibration error ?)

→ Zoom



T_e

He-data extrapolate well into the HCN & TS data



Summary & Conclusion

An extended model (to Brix 2000) with **proton collisions and CEX processes** has been tested on TEXTOR discharges.

Only marginal deviations => impact of the proton collisions is weak under the target plasma conditions investigated.

Evaluation with the non-stationary approach **(NS) enhances the radial extension** of the profile and **cancels relaxation phenomena**.

A method to apply the NS method as standard evaluation is under development.

For the measurement errors **the range of $\Delta T_e = 30\%$ and $\Delta n_e = 10\%$ given by Brix 2000 is still valid.**

Proton collisions (with $T_i = T_e$) do not exceed these margins.

Comparison of the $n_e(r)$ and $T_e(r)$ profiles obtained with other edge **are in reasonable agreement** with those from He.

BES on thermal He should be a reliable method for measurements within **$2.0 \times 10^{18} m^{-3} < n_e < 3.0 \times 10^{19} m^{-3}$ and $20 eV < T_e < 300 eV$.**

