

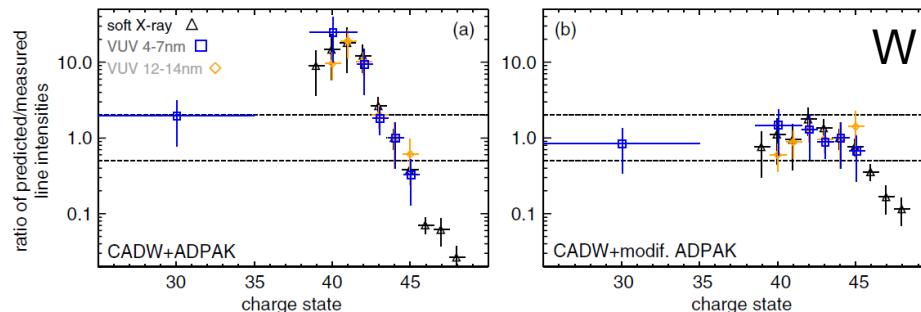
Role of atomic data in the ITC impurity transport code : the example of Ni

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CEA Cadarache

1. Principle : $n_z(z, r, t)$ calculation from hypothesis on transport
→ Role of ionisation / recombination rate coefficients
2. Method : match calculated emission with measurements
→ Role of emission coefficients
3. Consequences on determination of experimental transport

Introduction

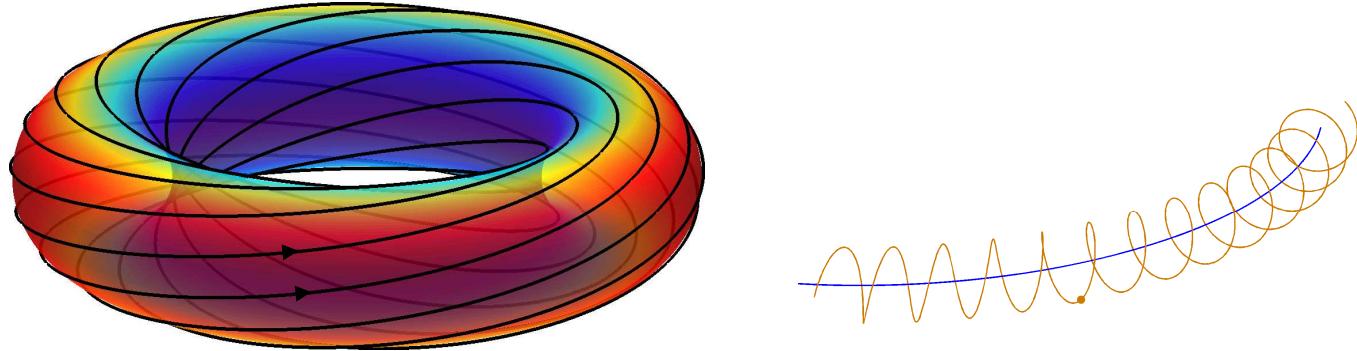
- Impurity transport codes (ITC, UTC, STRAHL,...) use many types of rate coefficients : ionisation, recombination, spectral line emission (UV, soft-X, CX), continuum
- Importance of data quality not easy to assess : measurements integrated over inhomogeneous plasma, rather lengthy minimisation process, many parameters
- It has been shown earlier that atomic data matter : e.g. Th. Pütterich PPCF 2008, but not much information available about effect on transport



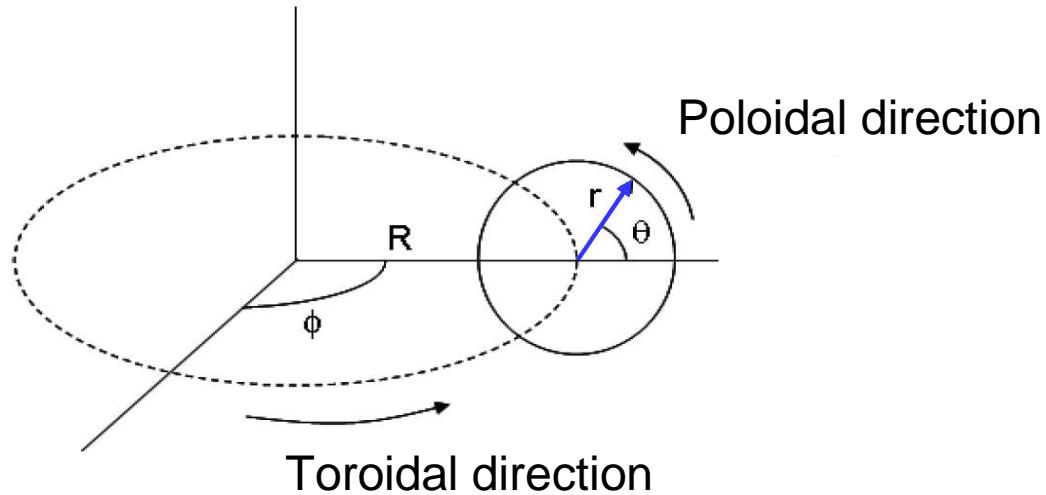
- Investigation made with ITC (CEA) : historical DB (Mattioli) / ADAS

(Tokamak configuration)

Magnetic field lines define closed, toric surfaces



Transport along field lines is very fast \Rightarrow assume toroidal, poloidal symmetry



Physics is assumed to depend only on the **minor radius r**

1. Principle : $n_z(z, r, t)$ calculation from hyp. transport

Background plasma ($n_e, T_e, \text{geometry}, \dots$)
Neutral source term



- Resolution of continuity equation for each ionisation stage

$$\frac{\partial n_z}{\partial t} + \nabla \Gamma_z = \alpha n_e n_z + S_{n_e n_z} + S_{Z_0}$$

**Recombination
and ionisation**

**Neutral
source**

- Transport model : **diffusion - convection**

$$\Gamma_z(r) = -D_z(r) \nabla_r n_z(r) + V_z(r) n_z(r)$$

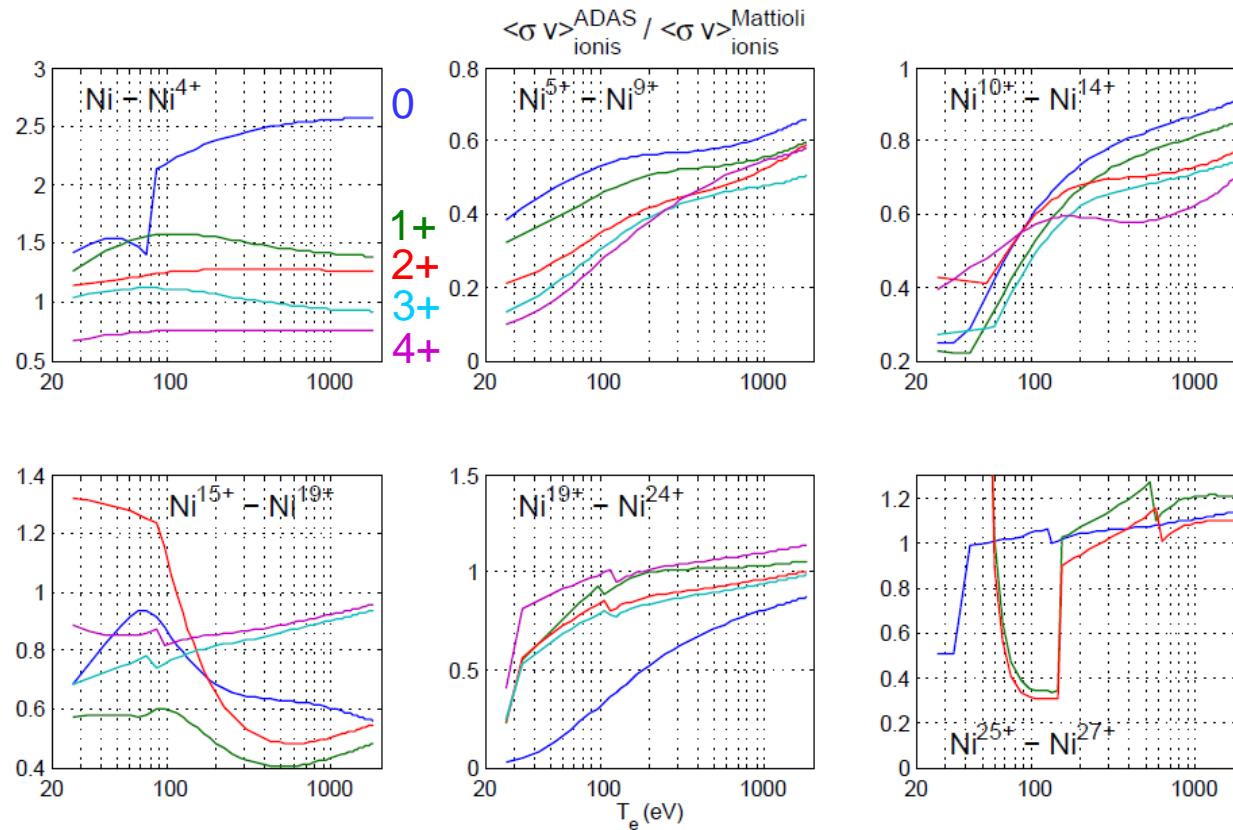
$$D_z = \begin{pmatrix} D_z & 0 & 0 \\ 0 & \ddots & 0 \\ 0 & 0 & D_z \end{pmatrix} \quad V_z = \begin{pmatrix} V_z & 0 & 0 \\ 0 & \ddots & 0 \\ 0 & 0 & V_z \end{pmatrix}$$

$$\alpha = \begin{pmatrix} 0 & \alpha_1 & 0 & & & 0 \\ 0 & -\alpha_1 & \alpha_2 & 0 & & 0 \\ 0 & 0 & -\alpha_2 & \ddots & 0 & 0 \\ & & 0 & \ddots & \ddots & \vdots \\ & & & 0 & -\alpha_{z-1} & \alpha_z \\ 0 & 0 & 0 & \dots & 0 & -\alpha_z \end{pmatrix} \quad \text{Recomb.}$$

$$S = \begin{pmatrix} -S_0 & 0 & 0 & & 0 \\ S_0 & -S_1 & 0 & & 0 \\ 0 & S_1 & -S_2 & 0 & 0 \\ 0 & \ddots & \ddots & 0 & \vdots \\ 0 & 0 & 0 & \ddots & -S_{z-1} & 0 \\ 0 & 0 & 0 & \dots & S_{z-1} & 0 \end{pmatrix} \quad \text{Ionis.}$$

1.1. Ionisation rate coefficients

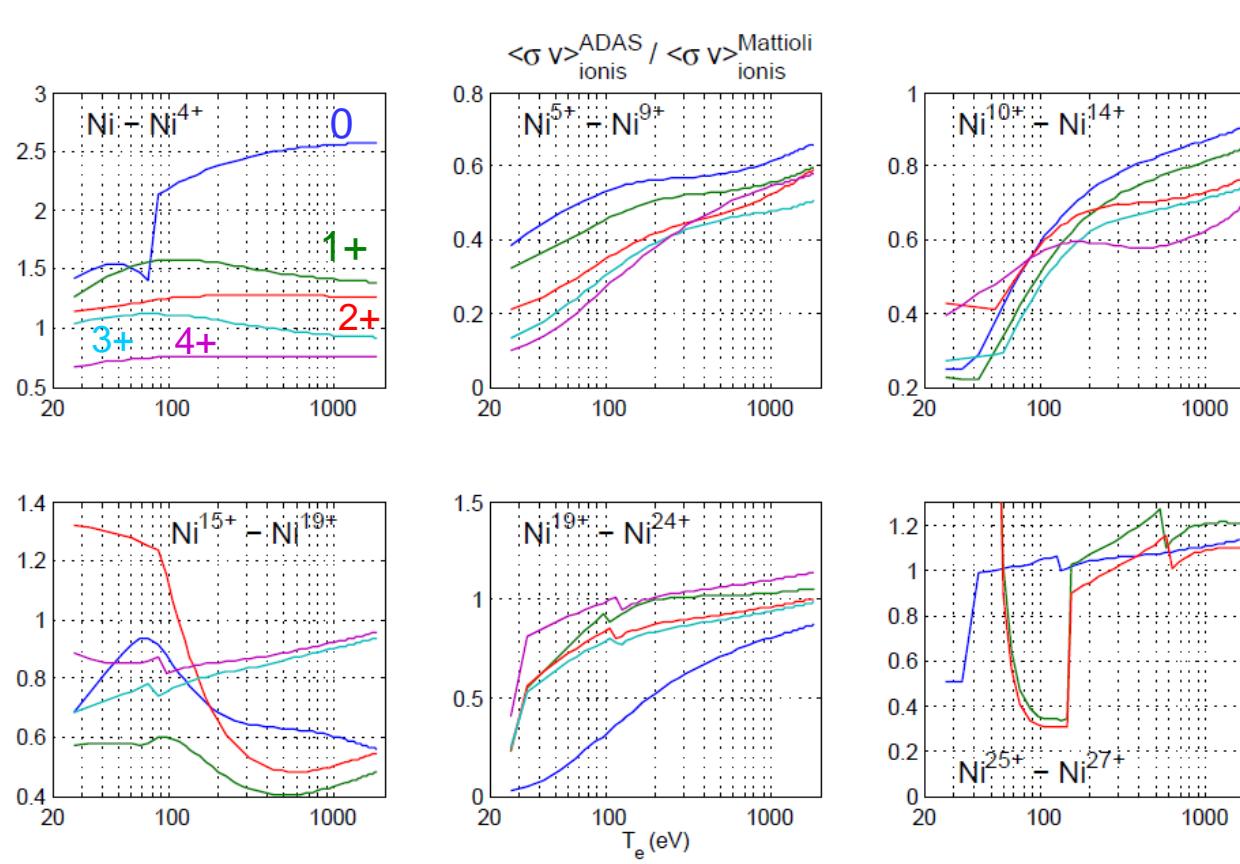
- **ADAS (solid)** : ADAS 408 on ADF03 atompars/atompars#ms_ni.dat
- **Mattioli (dashed)** : M.S. Pindzola et al., Physica Scripta T37, (1991) 35
(includes inner shell excit. + autoionis.)



- T_e dependences very similar
- Differences not much more than factor 5, gets smaller for higher ionisation stages

1.1. Ionisation rate coefficients

Ratios ADAS / Mattioli



Ratios ~ 1 for ions of greater importance (higher charge)

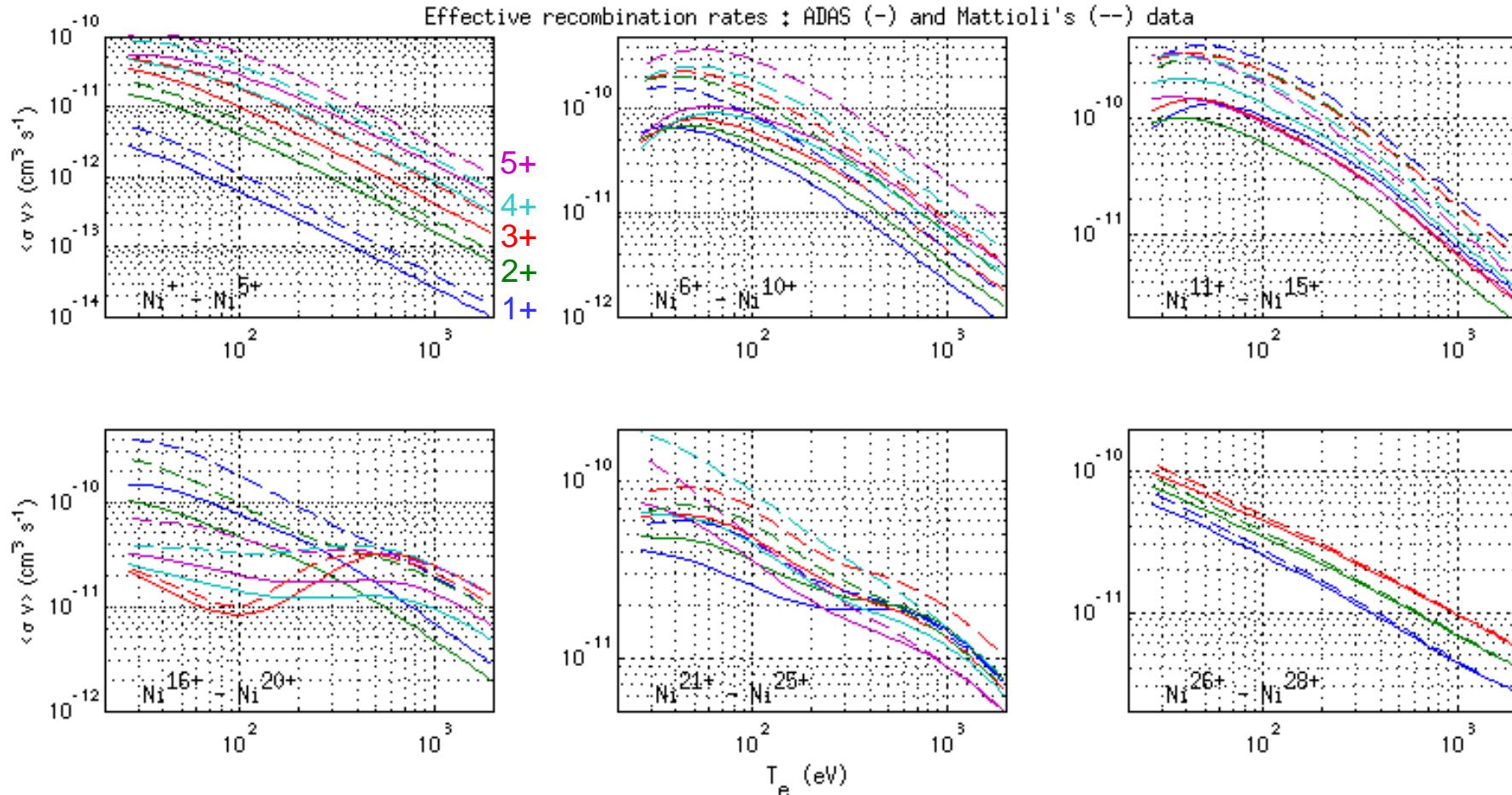
1.2. Recombination rate coefficients

- **ADAS (solid)** : ADAS408 on ADF03 atompars/atompars#ms_ni.dat:
Dielectronic : adjusted Burgess general formula or R-matrix (Badnell)
Radiative : analytic formula
- **Mattioli 1988 (dashed)** :
Dielectronic from various authors:
(H He) Arnaud & Rothenflug 1985 (Shull & Van Steenberg 1982)
(Li O F) Roszman 1987, (Be) Badnell 1987,
(B C N) Burgess 1965 - Merts 1976 + Badnell 1986
(Ne) Chen 1986 + Smith 1985
(Na Mg) Burgess 1965 - Merts 1976 + U. Connecticut 1984-7
(Al and >) Burgess 1965 - Merts 1976 ~ A-R

Radiative from formulae ~ Arnaud & Rothenflug 1985

1.2. Recombination rate coefficients

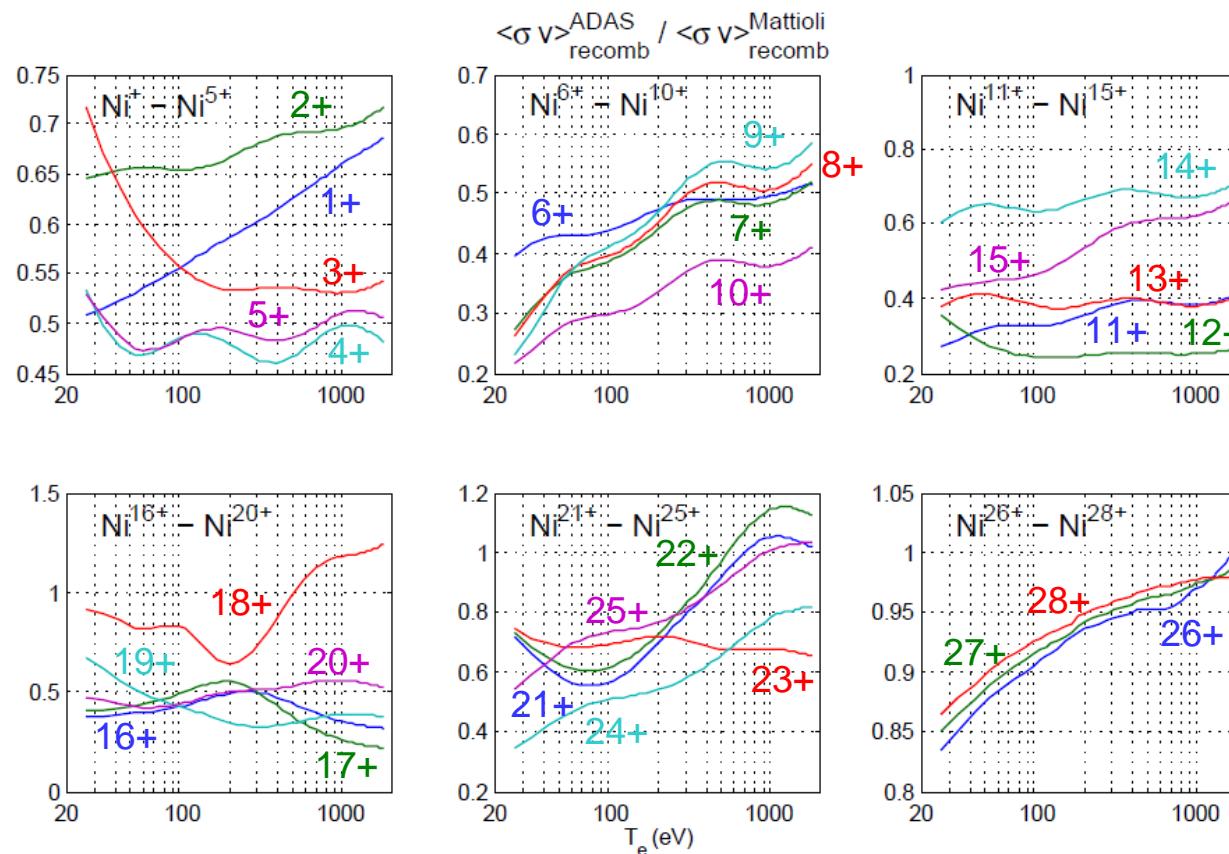
- ADAS (solid), Mattioli (dashed) :



- T_e dependences very similar
- Differences not much more than factor 2, very small for He-like to fully stripped

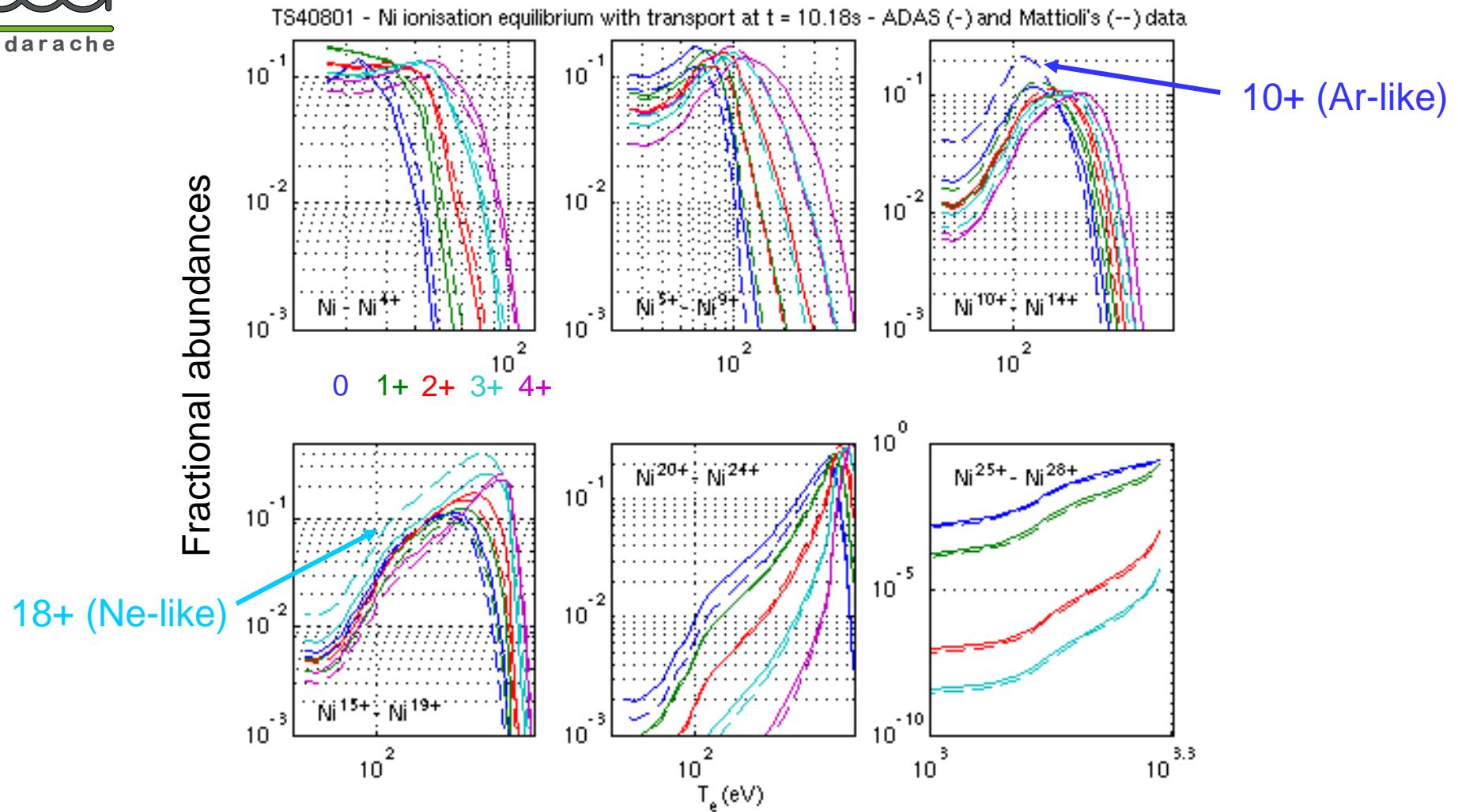
1.2. Recombination rate coefficients

Ratios ADAS / Mattioli



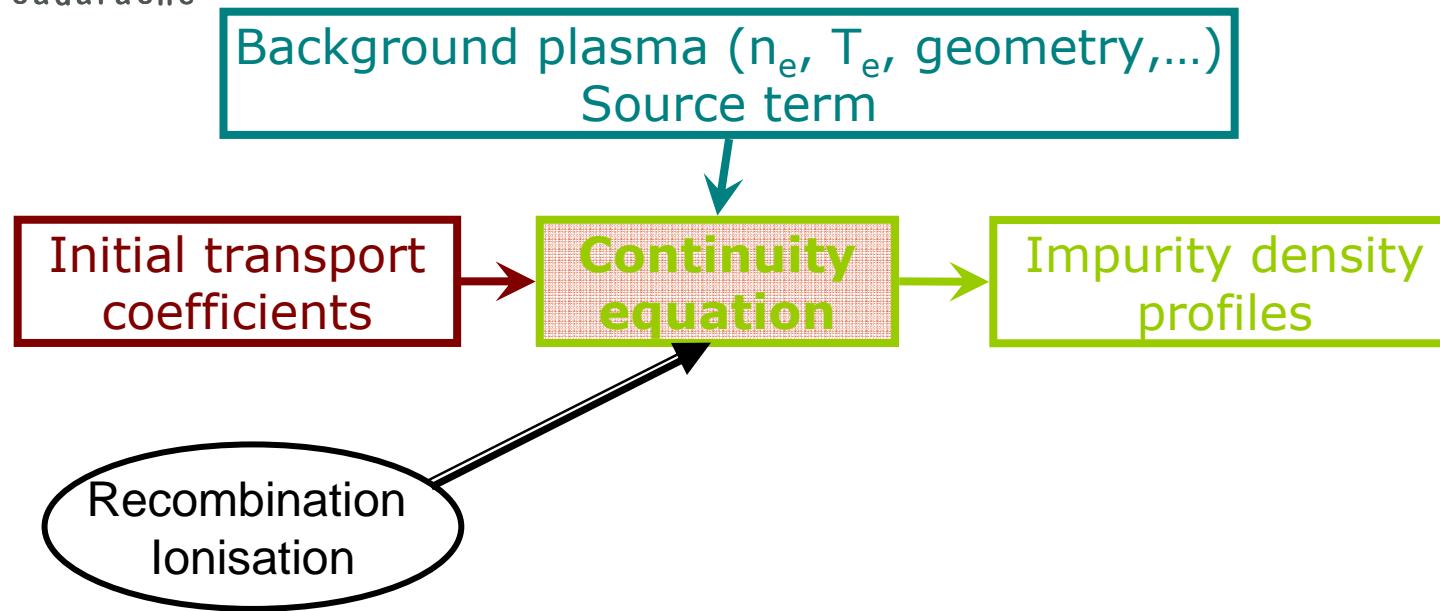
ADAS data always smaller than Mattioli but by a factor ≤ 2

1.3. Ionisation equilibrium (including transport)

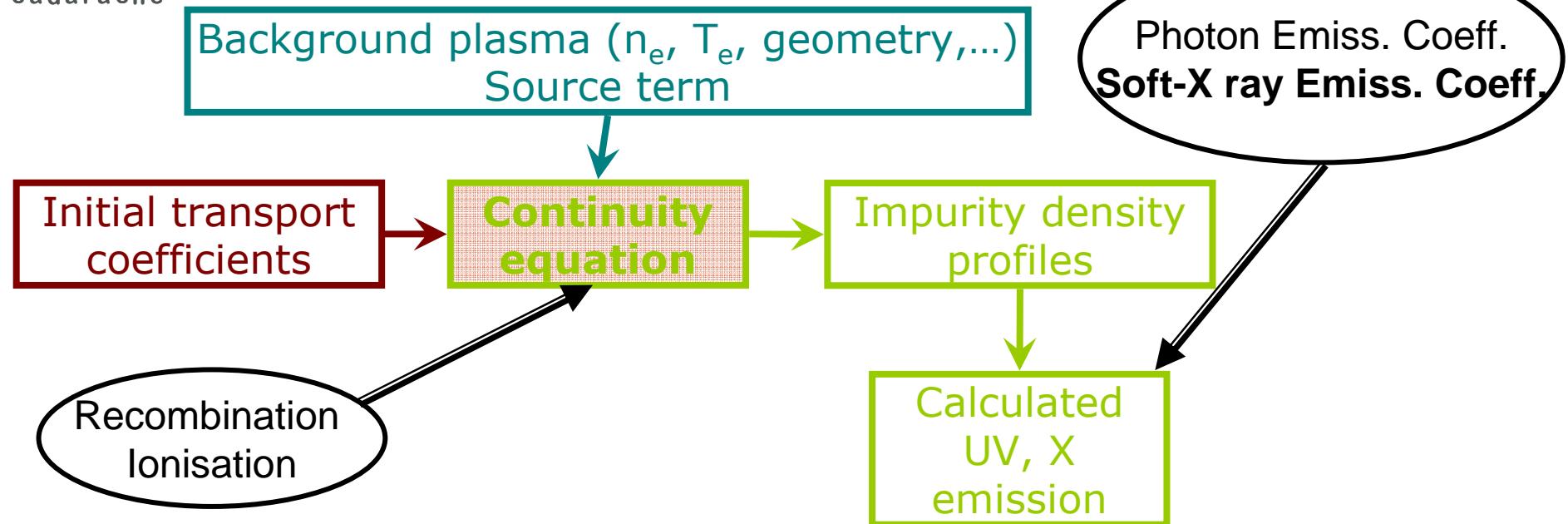


- Ionisation equilibrium not very different
- Ar-like and Ne-like, enhanced by 'Mattioli' data

2. Method : match calculated emission with measts

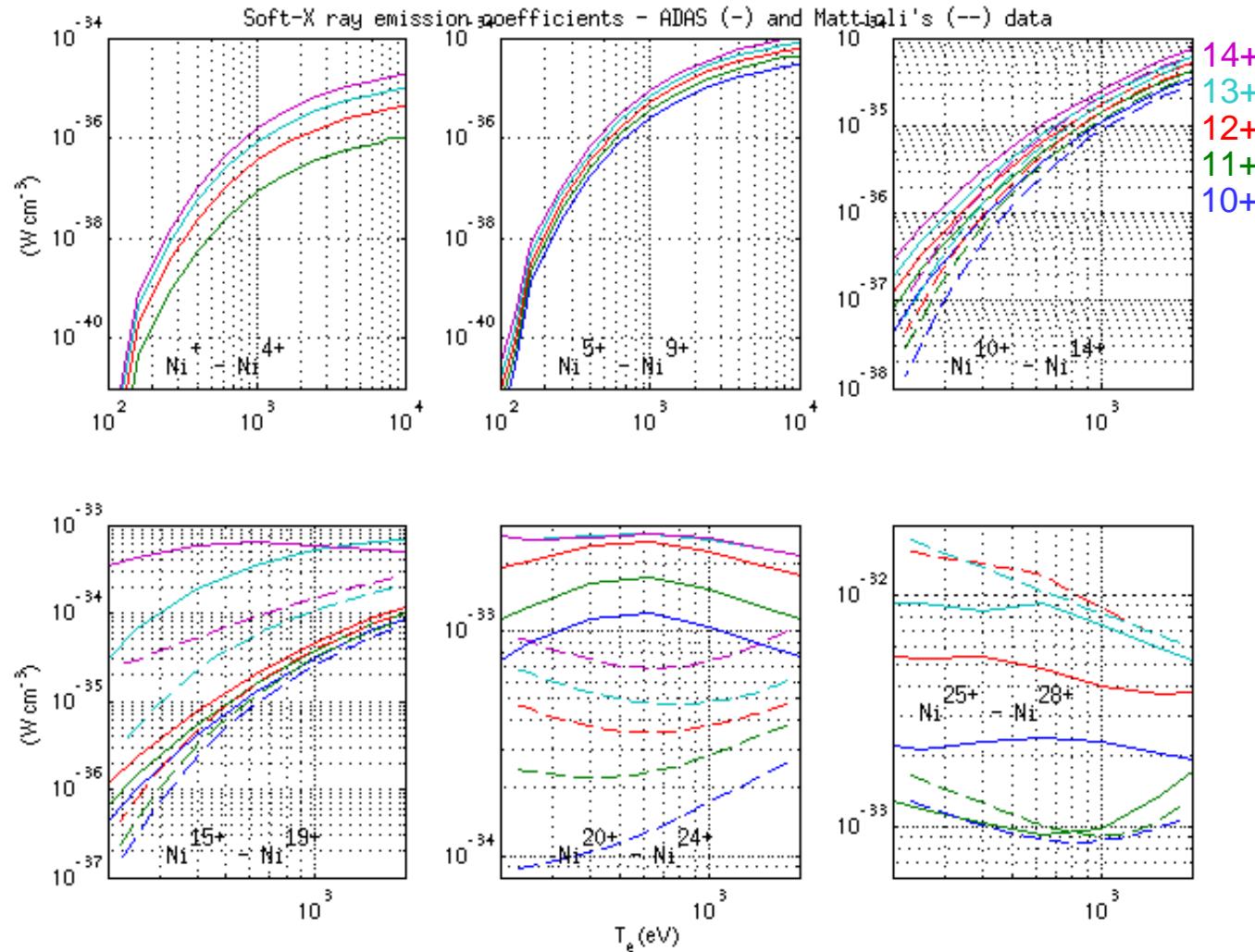


2. Method : match calculated emission with measts



2.1. Soft-X ray emission rate coefficients

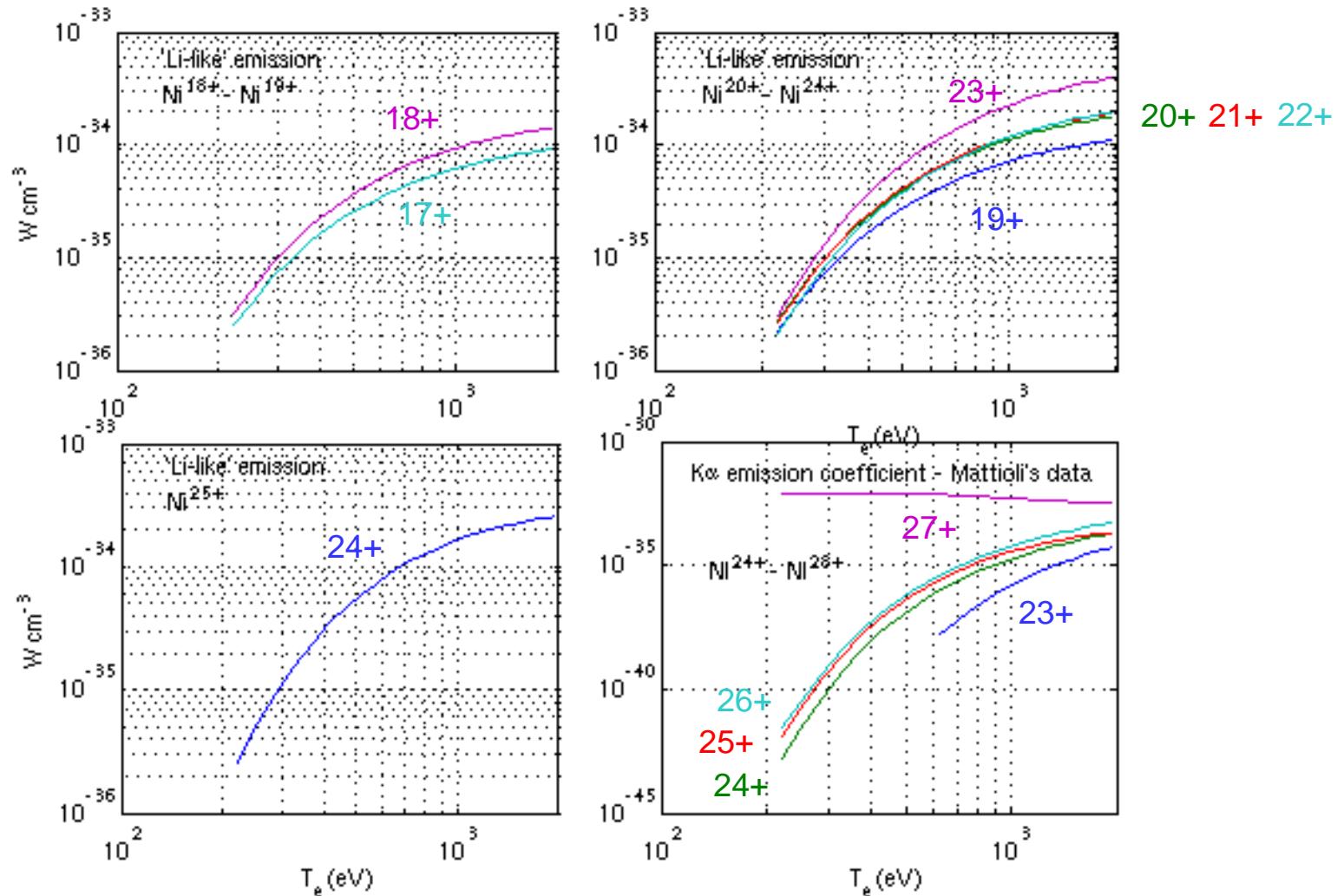
- ADAS (solid), Mattioli (dashed)



Large differences from 18+ (Ne-like) up

2.1. Soft-X ray emission rate coefficients

- Mattioli : separate K α and L spectrum contributions



2.2. Soft-X ray emission rate coefficients: detail

Mattioli :

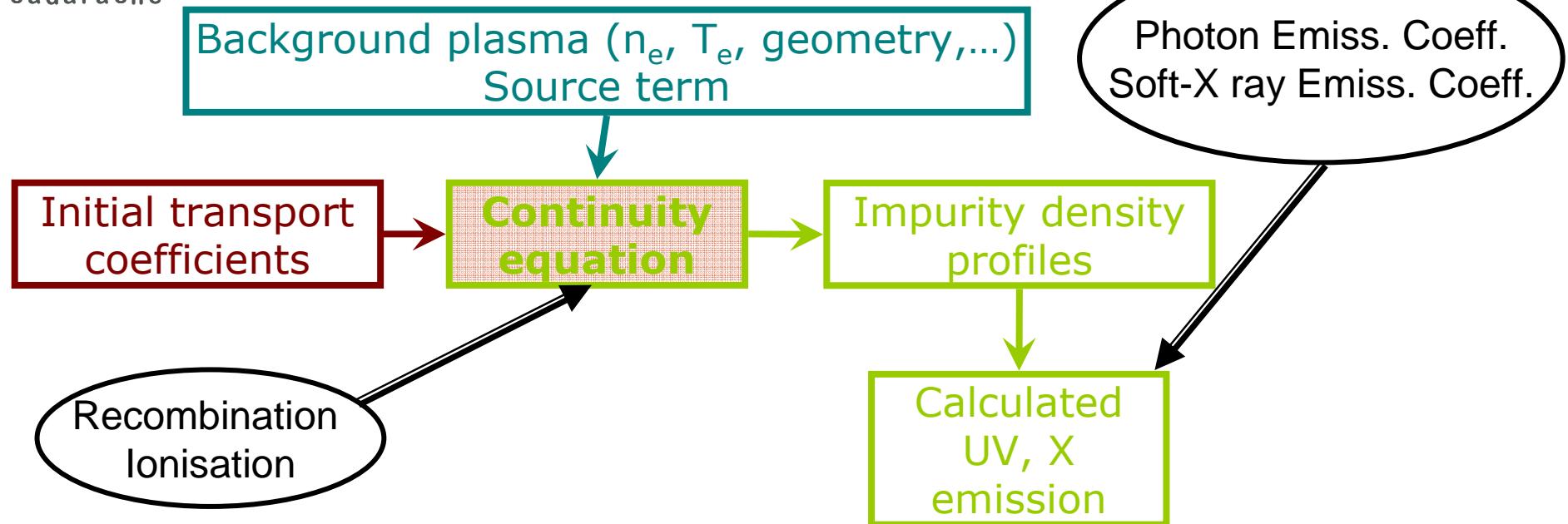
- Bremsstrahlung : known formulae + Hummer 1988, Carson 1987 (KL 1961)
Mattioli's code consistent with :

$$P_Z^{\text{Brems}} \propto \frac{n_e n_Z}{\sqrt{kT_e}} g_{ff} e^{-hv/kT_e} \quad (\text{proportionality coefficient?})$$

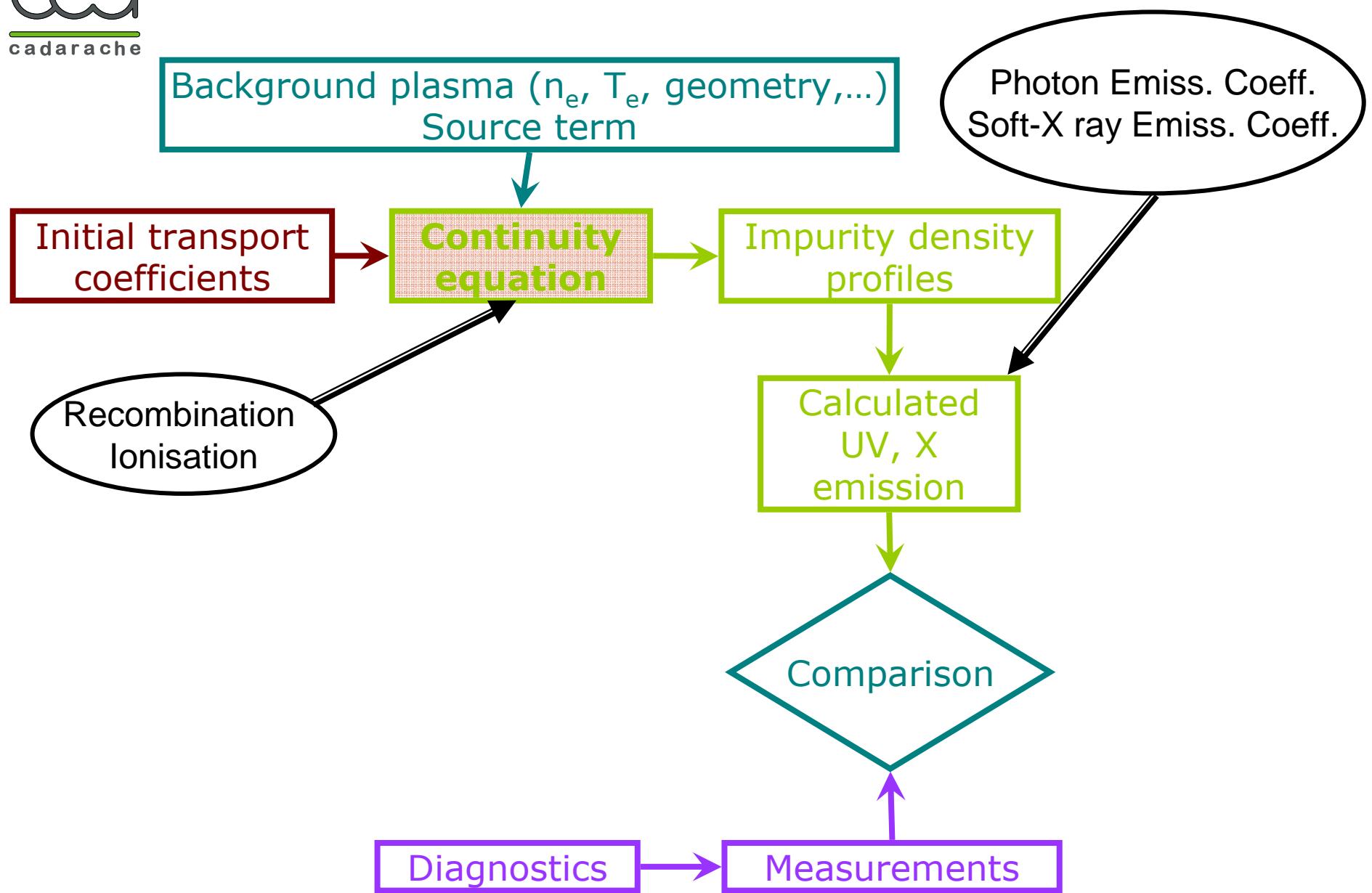
- Recombination: Burgess & Summers 1987 (KL 1961)
- K α : Clark 1982 H-like Ly α
 He-like w
Bombarda 1988 He-like (x+y+z)/w
 He-, Li-, Be-, B-like satellites
 H-like recombination contrib. to x, y, z, w
- L series: 2-3 transitions for Li- to Ne-like \rightarrow analytic formula
B-, C-, N- and Ne-like : Bhatia 1985-89
Li-like : Cochrane 1983

ADAS :

2. Method : match calculated emission with measts

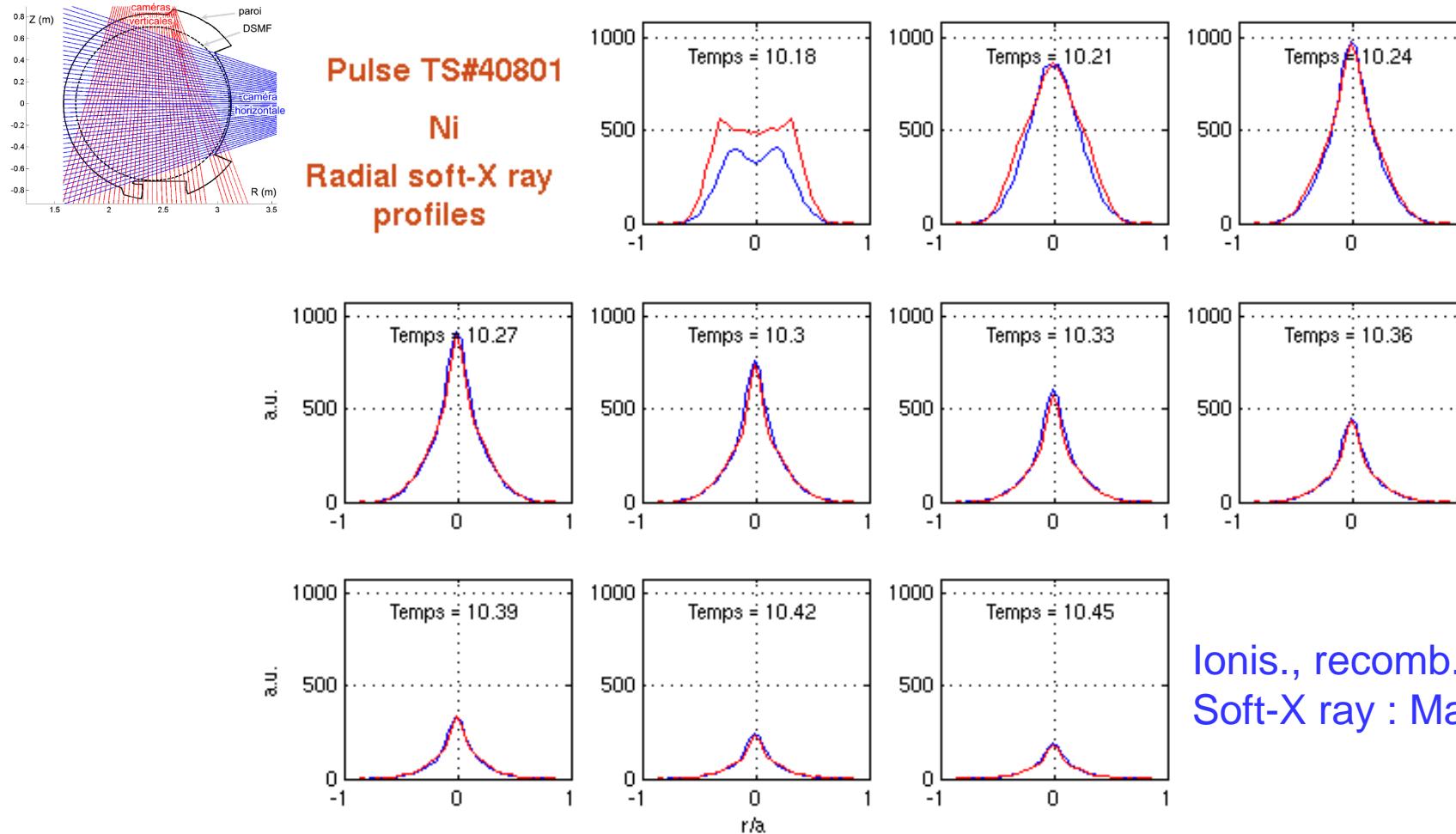


2. Method : match calculated emission with measts



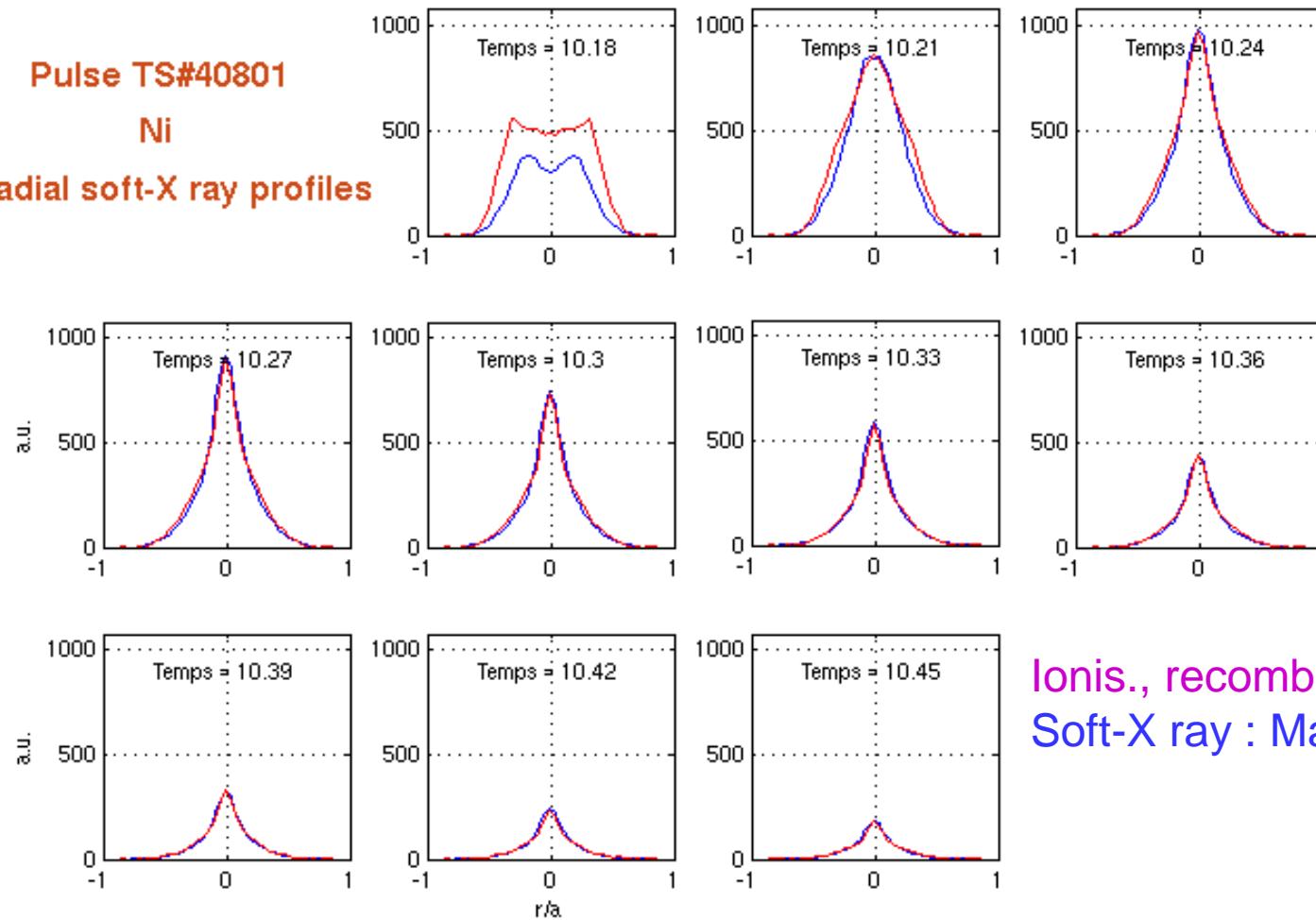
3.1. Effect of ionis./recomb. coefficients in emission calculations

- Optimisation of the transport coefficients has been performed with 'Mattioli' data



3.1. Effect of ionis./recomb. coefficients in emission calculations

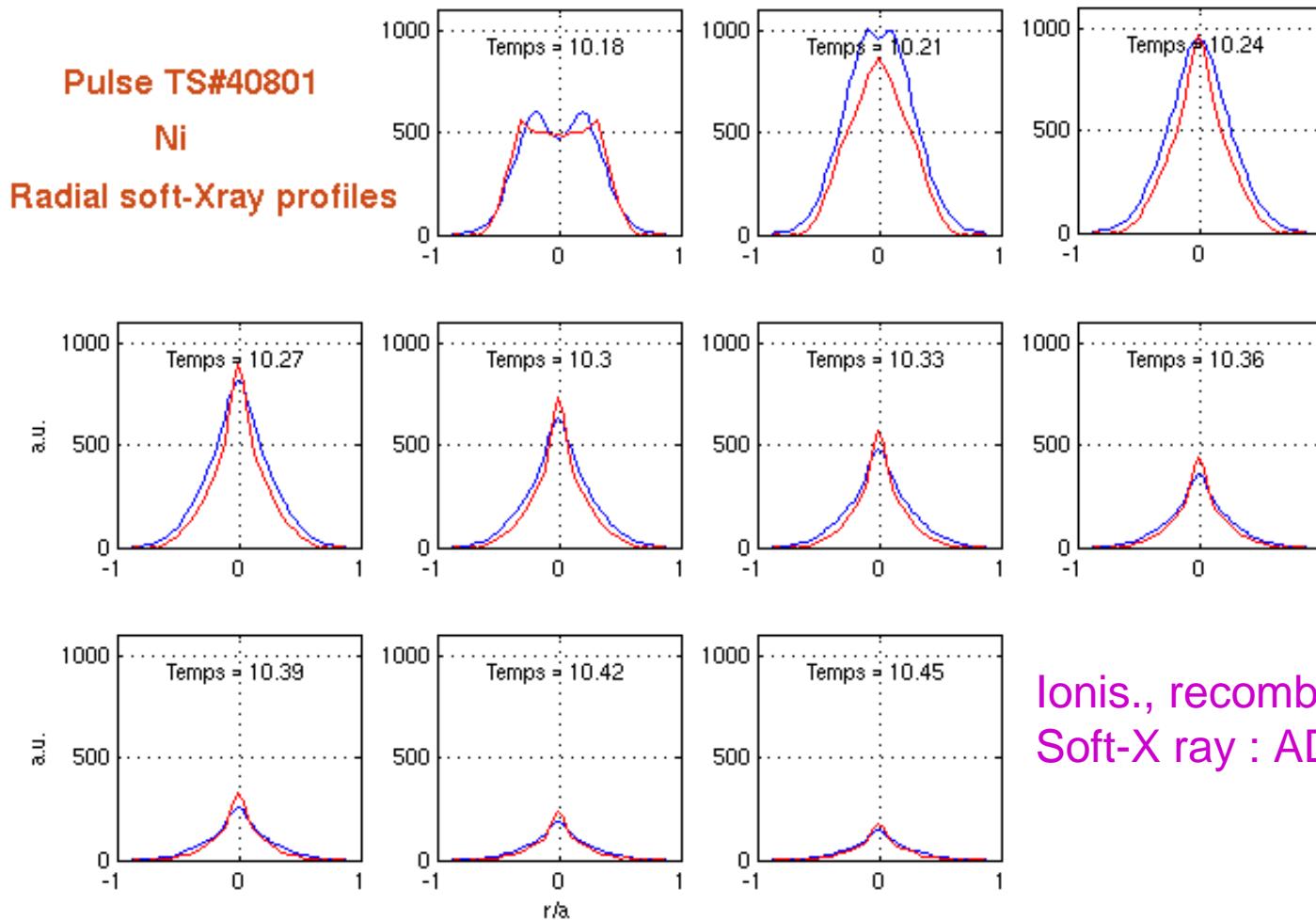
Pulse TS#40801
Ni
 Radial soft-X ray profiles



Ionis., recomb. : ADAS
 Soft-X ray : Mattioli

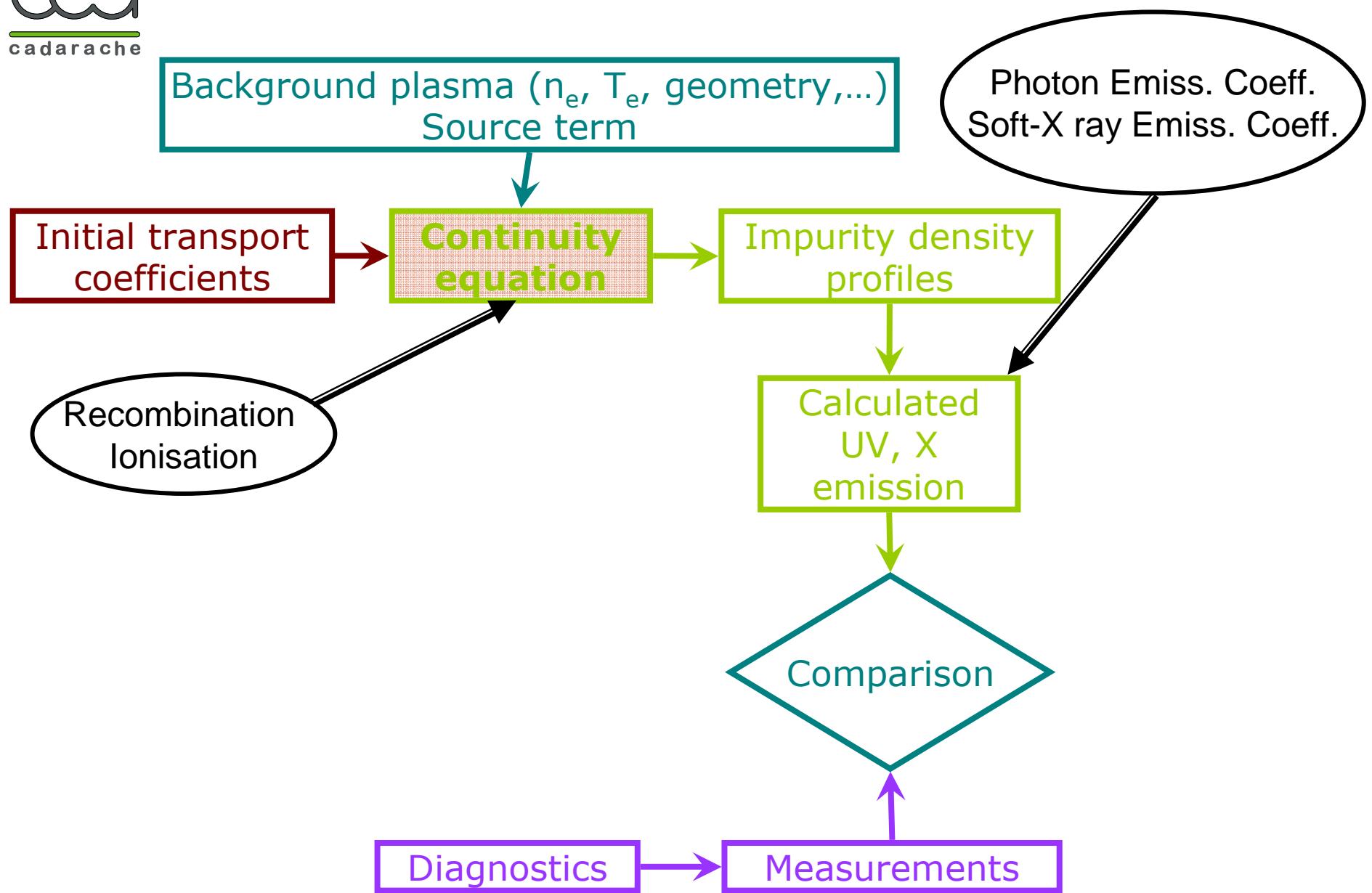
Calculated emission weakly sensitive to ionisation equilibrium
 (coefficients within factor 2)

3.2. Effect of soft-X ray coefficients in emission calculations

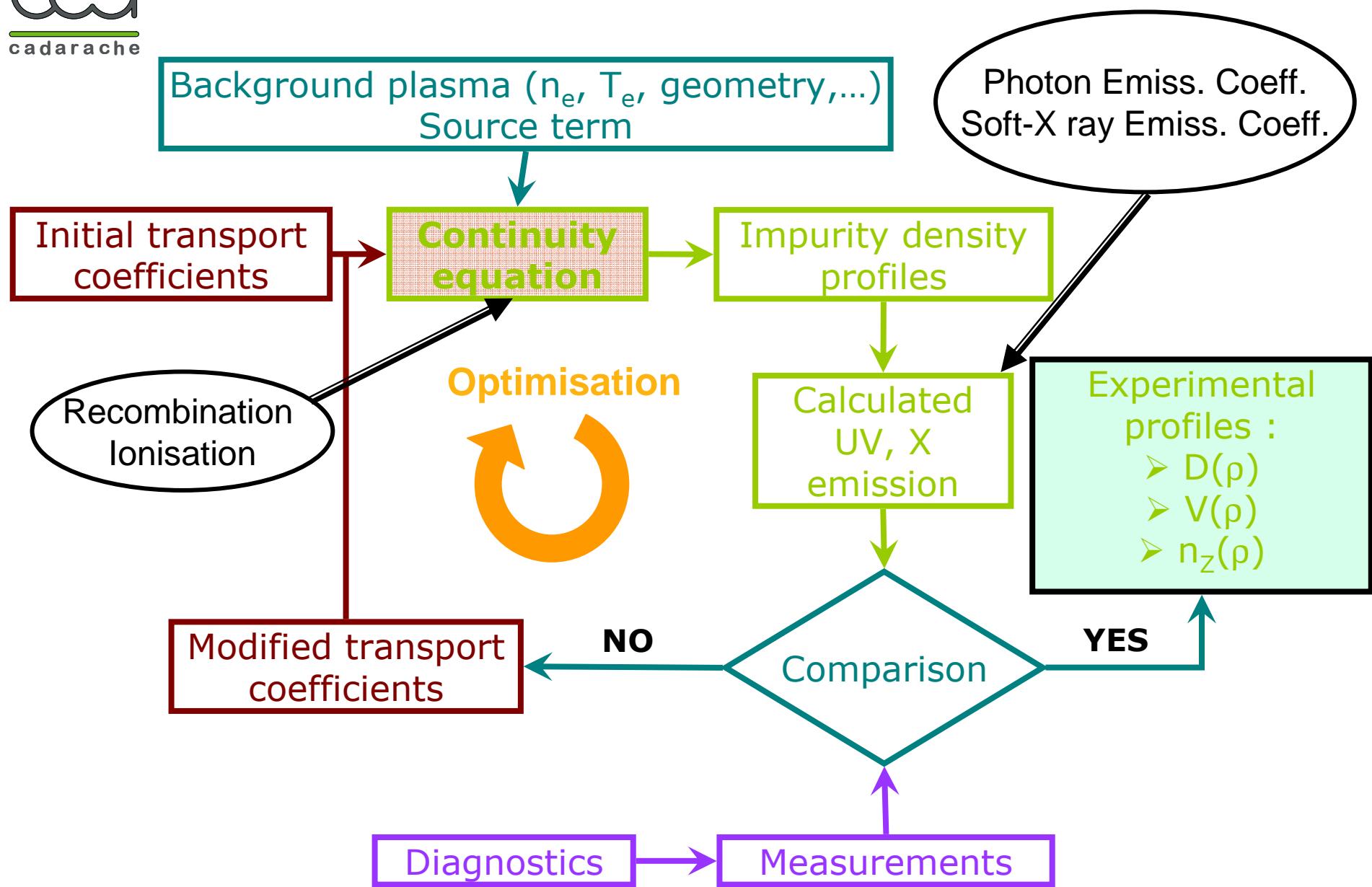


Calculated emission very sensitive to soft-X ray emission coefficients

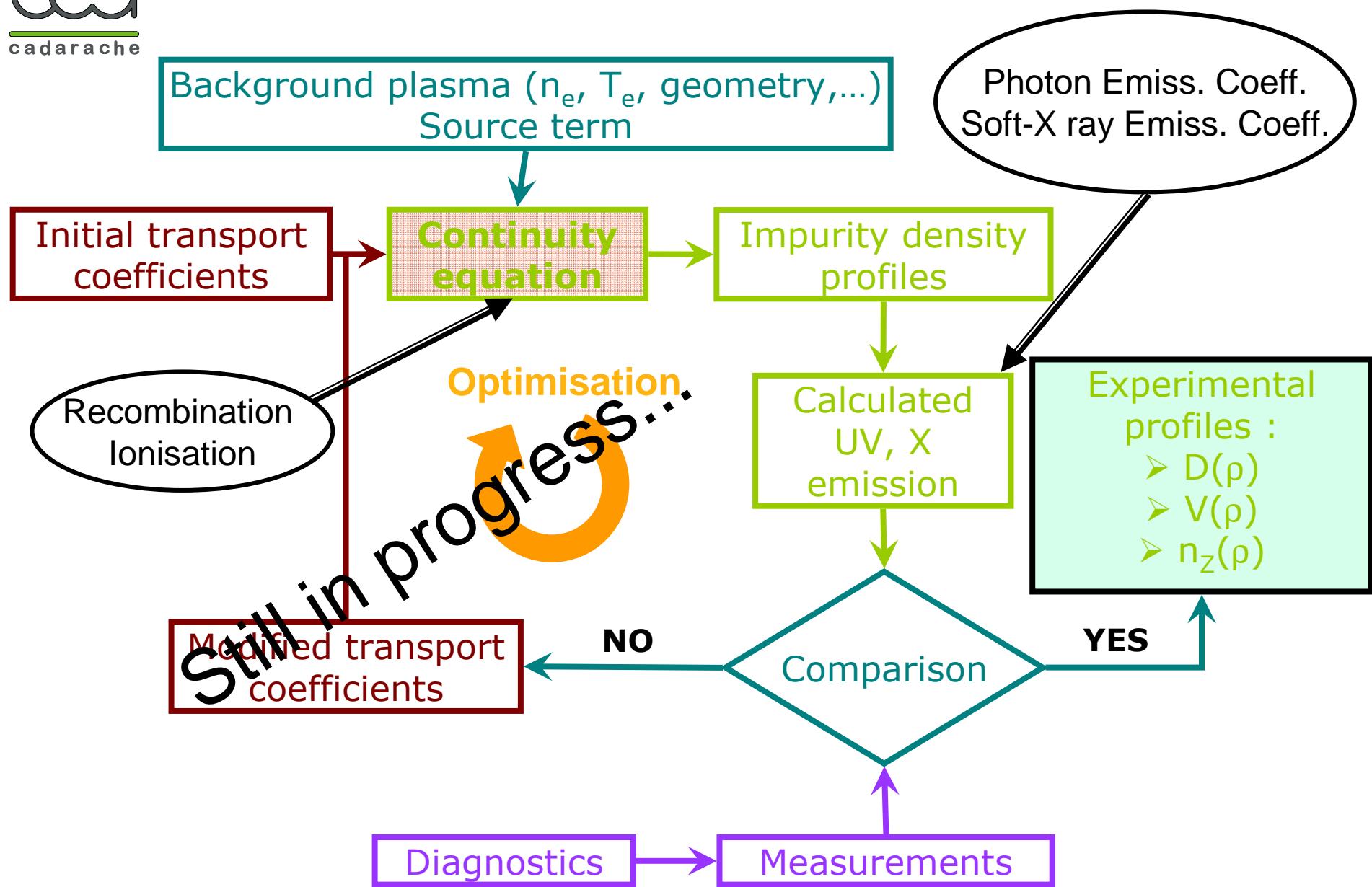
2. Method : match calculated emission with measts



2. Method : match calculated emission with measts



2. Method : match calculated emission with measts



Conclusions

Atomic physics:

- Ionisation

Factors up to 5 between ADAS and ‘Mattioli’ rate coefficients

Difference smaller for higher charge states

- Recombination

ADAS rate coeff. smaller by a factor 1-3

→ Ionisation equilibrium : Mattioli’s coefficients favour Ar- and Ne-like

- **Soft-X ray emiss. coefficients: large differences (~ 10) for Ne-like and above**

Transport:

- **Soft-X ray emission coefficients essential** (value + T_e dependence)

- Small differences in ionisation equilibrium have a weak effect on calculated emission (soft-X ray, UV)

- UV PEC T_e dependence relatively unimportant for emitters confined on narrow layers, absolute values must be consistent with soft-X ray coeff.

- In progress : propagation of atomic physics uncertainties to transport

- Ionisation rate coefficients

M.S. Pindzola et al., Physica Scripta T37, (1991) 35 (*includes inner shell excit. + autoionis.*)

- Recombination rate coefficients

→ *Dielectronic from various authors:*

- (H He) Arnaud & Rothenflug, Astron. Astrophys Suppl.Ser. 60 (1985) 425
Schull & Van Steenberg, Astrophys. Suppl. Ser. 48 (1982) 95
- (Li O F) Roszman, Phys. Rev. A 35 (1987) 2122
- (Be) Badnell, J. Phys. B 20 (1987) 2081
- (B C N) Burgess, Astrophys. J. 141 (1965) 1588
Merts, Los Alamos Scientific Lab. Report LA-6220-MS (1976)
Badnell, J. Phys. B 20 (1986) 3827
Roszman, Phys. Rev.A 35 (1987) 2138 & 3368
- (Ne) Chen M.H., Phys. Rev. A 34 (1986) 994
Smith B.W., Astrophys. J. 298 (1985) 898
- (Na Mg) Burgess – Merts
Lagatutta, Phys. Rev. A 30 (1984) 316
Dube, J. Quant. Spectrosc. Radiat. Transfer 33 (1985) 13
Dube, J. Quant. Spectrosc. Radiat. Transfer 38 (1987) 311
- (Al and >) Burgess - Merts ~ A-R

→ *Radiative from formulae ~ A-R*

→ Radiative from formulae ~ A-R

Fully stripped : $\alpha_{rZ} = 5.2 \times 10^{-14} Z^2 \sqrt{\frac{I_H}{T_e}} \phi_1(\beta)$ with $I_H = 13.6 eV$

$$\beta = Z^2 \frac{I_H}{T_e}$$

$$\phi_1(\beta) = \sum_{n=1}^{\infty} \frac{\beta}{n^3} e^{\beta/n^2} E_1\left(\frac{\beta}{n^2}\right)$$

Others : $\alpha_{rZ} = 2.6 \times 10^{-14} (\alpha_1 + \alpha_2)$ with

$$\alpha_1 = Z^2 \sqrt{\frac{I_H}{T_e}} \frac{\xi}{n^3} \frac{I_{Z-1}}{T_e} e^{I_{Z-1}/T_e} E_1\left(\frac{I_{Z-1}}{T_e}\right)$$

$$\alpha_2 = 2Z^2 \sqrt{\frac{I_H}{T_e}} \phi_{n+1}\left(Z^2 \frac{I_H}{T_e}\right)$$

(ξ empty sites in valence shell n)

- Soft-X ray emission rate coefficients
 - Bremsstrahlung: known formulae
 - Hummer, *Astrophys. J.* 327 (1988) 477
 - Karzas, *Astrophys. J. Suppl. Ser.* 6 (1961) 167
 - Carson, *Astronom. Astrophys.* 189 (1988) 319
 - Recombination: Burgess & Summers, *Mon. Not. R. Astr. Soc.* 226 (1987) 257
 - K α spectrum: Clark, *Astrophys. J.* 254 (1982) 412
 - H-like Ly α
 - He-like w
 - Bombarda, *Phys. Rev. A* 37 (1988) 504
 - He-like (x+y+z)/w
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 - L spectrum: 2-3 transitions for Li- to Ne-like → analytic formula
 - B-, C-, N- and Ne-like : Bhatia, *ADNDT* 32 (1985) 435, 35 (1986) 319, 35 (1986) 449, 36 (1987) 453, 43 (1989) 99
 - Li-like : Cochrane, *Physica Scripta* 28 (1983) 25

Th. Puetterich, PPCF 2008

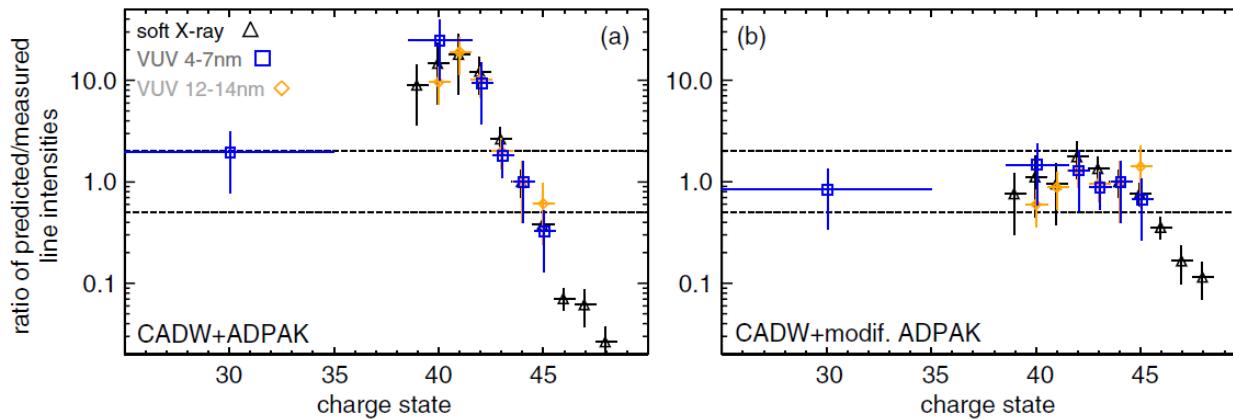


Figure 5. (a) Ratios of predicted to measured line intensities for different wavelength ranges versus ionization stage using the ‘CADW+ADPAK’ ion balance (see figure 6). Several lines of an ionization stage are summed in the model and in the spectrum to reduce the uncertainties that could occur for a single spectral line. (b) Similar data as (a), but using the data set ‘CADW+modif. ADPAK’ which is described in the text. Dashed lines correspond to ‘factor of 2’-margin around 1 introduced to guide the eye.